

ON A QUASI-NORMAL  $(\psi, g, u, v, e, \lambda)$ -STRUCTURE

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In the present paper it is shown that the  $(\psi, g, u, v, e, \lambda)$ -structure is quasi-normal in the sense of Yano (Yano and Ki 1972) iff the vector fields are harmonic. Also, the induced structure of the invariant submanifold of co-dimension 2 is quasi-normal.

1. GENERAL FORMULA FOR  $(\psi, g, u, v, e, \lambda)$ -STRUCTURE

Let  $V_n$  be a  $C^\infty$  differentiable manifold with (1, 1) tensor field  $\psi$ , Riemannian metric tensor  $g$ , three 1-forms  $u, v, e$  and associated with three vector fields  $U, V, E$  and a differentiable function  $\lambda$  satisfying (Baik 1972, Yano 1965)

- (a)  $\psi_j^i \psi_i^h = -\delta_j^h + u_j u^h + v_j v^h + e_j e^h,$
- (b)  $\psi_j^i \psi_i^s g_{ts} = g_{ji} - u_j u_i - v_j v_i - e_j e_i,$
- (c)  $\psi_i^t u_t = \lambda v_i, \psi_i^t v_t = -\lambda u_i, \psi_i^t e_t = 0,$
- (d)  $\psi_i^h u^i = -\lambda v^h, \psi_i^h v^i = \lambda u^h, \psi_i^h e^i = 0,$
- (e)  $u_i u^i = 1 - \lambda^2, u_i v^i = 0, u_i e^i = 0,$
- (f)  $v_i u^i = 0, v_i v^i = 1 - \lambda^2, v_i e^i = 0,$
- (g)  $e_i u^i = 0, e_i v^i = 0, e_i e^i = 1.$  ...(1.1)

Let us put

$$\psi_{ji} = \psi_j^t g_{ti} \quad (\text{being skew symmetric}) \tag{1.2}$$

$h, i, j, \dots$  run over  $1, 2, \dots, n = (2m + 3).$

Let us define (1.2) type of torsion tensor as follows:

$$S_{ji}^{def} = \psi_j^i \nabla_i \psi_i^h - \psi_i^j \nabla_j \psi_i^h - \psi_i^h (\nabla_j \psi_i^t - \nabla_i \psi_j^t) + u_j u^h + v_j v^h + e_j e^h \tag{1.3}$$

where  $u_i = \nabla_j u_i - \nabla_i u_j$  is covariant differential w.r.t. Levi Civita connection. If the torsion tensor  $S_{ji}^h$  vanishes, the  $(\psi, g, u, v, e, \lambda)$ -structure is called normal.

Transvecting (1.3) with  $u_h, v_h, e_h$  and also keeping in view (1.1) we obtain

$$S_{ji}^h u_h = u_{ji} + \lambda(u_i \nabla_j \lambda - u_j \nabla_i \lambda) + \lambda(\psi_j^t v_{ii} - \psi_i^t v_{ij}) + \nabla_i \lambda(\psi_j^t v_i - \psi_i^t v_j) - \psi_i^s \psi_j^s u_{is}, \tag{1.4}$$

$$S_{ji}^h v_h = v_{ji} + \lambda(v_i \nabla_j \lambda - v_j \nabla_i \lambda) - \lambda(\psi_j^t u_{ii} - \psi_i^t u_{ij}) - \nabla_i \lambda(\psi_j^t u_i - \psi_i^t u_j) - \psi_j^t \psi_i^s v_{is}, \tag{1.5}$$

and

$$S_{ji}^h e_h = e_{ji} + \psi_j^t \psi_i^s e_{si}. \tag{1.6}$$

We have (Yano and Ki 1972)

$$\psi_{jih} = \nabla_j \psi_{ih} + \nabla_i \psi_{hj} + \nabla_h \psi_{ji}. \tag{1.7}$$

Using (1.3) and (1.7) the covariant components  $S_{jih}$  can be written as

$$S_{jih} - (\psi_j^t \psi_{iih} - \psi_i^t \psi_{ijh}) = \psi_i^t \nabla_h \psi_{ij} - \psi_j^t \nabla_h \psi_{ii} + u_i \nabla_i u_h + v_j \nabla_i v_h + e_j \nabla_i e_h - u_i \nabla_j u_h - v_i \nabla_j v_h - e_i \nabla_j e_h. \tag{1.8}$$

Transvecting (1.8), with  $u^j, v^j, e^j$ , we get respectively

$$\{S_{jih} - (\psi_j^t \psi_{iih} - \psi_i^t \psi_{ijh})\} u^j = -\lambda^2 u_{ih} + \underset{u}{l} g_{ih} - u_i u^t \underset{u}{l} g_{ih} - e_i u^t e_{ih} + \lambda \psi_i^t \underset{v}{l} g_{ih} - v_{ih} (v_i u^t + \lambda \psi_i^t) \tag{1.9}$$

$$\{S_{jih} - (\psi_j^t \psi_{iih} - \psi_i^t \psi_{ijh})\} v^j = -\lambda^2 v_{ih} + \underset{v}{l} g_{ih} - v_i v^t \underset{v}{l} g_{ih} - e_i v^t e_{ih} - \lambda \psi_i^t \underset{u}{l} g_{ih} + (\lambda \psi_i^t - u_i v^t) u_{ih} \tag{1.10}$$

and

$$\{S_{jih} - (\psi_j^t \psi_{iih} - \psi_i^t \psi_{ijh})\} e^j = \underset{e}{l} g_{ih} - e_i e^t \underset{e}{l} g_{ih} - v_i e^t v_{ih} - u_i e^t u_{ih}. \tag{1.11}$$

where  $\underset{u}{l}, \underset{v}{l}, \underset{e}{l}$  denote the Lie derivative w.r.t. the vector field  $u, v, e$  respectively.

### 2. QUASI-NORMAL $(\psi, g, u, v, e, \lambda)$ -STRUCTURE

$(\psi, g, u, v, e, \lambda)$ -structure is called quasi-normal if (Yano and Ki 1972)

$$S_{jih} - (\psi_j^t \psi_{iih} - \psi_i^t \psi_{ijh}) = 0. \tag{2.1}$$

Using (2.1), eqns. (1.9) – (1.11) reduce to

$$\begin{aligned} l g_{ih} - u_i u^t l g_{th} + \lambda \psi_i^t l g_{th} - e_i u^t e_{th} \\ = \lambda^2 u_{ih} + v_{ih}(v_i u^t + \lambda \psi_i^t) \end{aligned} \quad \dots(2.2)$$

$$\begin{aligned} l g_{ih} - v_i v^t l g_{th} - \lambda \psi_i^t l g_{th} - e_i v^t e_{th} \\ = \lambda^2 v_{ih} - u_{ih}(\lambda \psi_i^t - u_i v^t) \end{aligned} \quad \dots(2.3)$$

and

$$l g_{ih} - e_i e^t l g_{th} - v_i e^t v_{th} - u_i e^t u_{th} = 0. \quad \dots(2.4)$$

Substituting (2.2) and (2.3) in (2.4) and on taking account of (1.1), we obtain

$$\lambda^2 l g_{ih} - \lambda^2 e_i e^t l g_{th} + v_i v^s e_{sh} - v_i e^t l g_{th} - u_i e^t l g_{th} + u_i u^s e_{sh} = 0. \quad \dots(2.5)$$

Transvecting (2.5) with  $u^i$  and using (1.1), we get

$$u^s e_{sh} = e^t l g_{th} - \frac{\lambda^2}{1 - \lambda^2} u^i l g_{ih}. \quad \dots(2.6)$$

Substituting the value of  $u^s e_{sh}$  in eqn. (2.5) and then transvecting the resulting equation with  $e^i$  and  $v^i$  respectively, we find

$$l g_{ih} = 0 \quad \dots(2.7a)$$

and

$$v^s e_{sh} = e^t l g_{th}. \quad \dots(2.7b)$$

From (1.1), (2.2) and (2.3), we have

$$\begin{aligned} \lambda^2(1 - \lambda^2) v_{ih} + v_{ih} \{u_i u^s - \lambda^2 v_i v^s - \lambda^2 e_i e^s\} + \{\lambda(1 - \lambda^2) \psi_i^t \\ - \lambda^2 v_i u^t - u_i v^t\} l g_{th} + \lambda^2 e_i e^t l g_{th} - \lambda^2 e_i v^t e_{th} = 0 \end{aligned} \quad \dots(2.8)$$

which on transvecting with  $u^i$  and in view of (1.1) becomes

$$u^s v_{sh} = v^s l g_{sh} \quad (\lambda \neq \pm 1, \pm i) \quad \dots(2.9)$$

From (2.9), it can be seen that

$$(l g_{is}) u^j v^t = 0. \quad \dots(2.10)$$

Transvecting (2.2) and (2.3) with  $e^t$ , we get

$$e^t \underset{u}{l} g_{ih} - u^t e_{ih} = \lambda^2 u_{ih} e^t \quad \dots(2.11a)$$

and

$$e^t \underset{v}{l} g_{ih} - v^t e_{ih} = \lambda^2 e^i v_{ih}. \quad \dots(2.11b)$$

Transvecting (2.2) with  $v^i$  and also using (1.1) and (2.9), we find

$$u^t \underset{u}{l} g_{ih} = v^t u_{hi}. \quad \dots(2.12)$$

Making use of (2.6) and (2.7a, b) in (2.11a, b) we obtain

$$(a) \quad u_{ih} = 0, \quad (b) \quad v_{ih} = 0, \quad (c) \quad e_{ih} = 0. \quad \dots(2.13)$$

Similarly from (2.9), (2.12) and (2.13), we have

$$(a) \quad \underset{u}{l} g_{ih} = 0, \quad \text{and} \quad (b) \quad \underset{v}{l} g_{ih} = 0. \quad \dots(2.14)$$

Since  $\underset{u}{l} g_{sh} = \nabla_s u_h + \nabla_h u_s = 0$ .

which can be written as

$$\nabla_s u^s = 0.$$

Thus in a quasi-normal  $(\psi, g, u, v, e, \lambda)$ -structure, we get  $u_{ih} = 0, v_{ih} = 0, e_{ih} = 0$  and  $\nabla_i u^i = 0, \nabla_i v^i = 0$  and  $\nabla_i e^i = 0$ .

Now if we substitute the last set of equations in (1.9) – (1.11), we find

$$S_{jih} - (\psi_j^i \psi_{ih} - \psi_i^j \psi_{jh}) = 0$$

which is the condition for a structure to be quasi-normal. Hence

*Theorem 2.1* — If the  $(\psi, g, u, v, e, \lambda)$ -structure is quasi-normal, it is necessary and sufficient that  $U, V, E$  are harmonic vectors ( $\lambda \neq \pm 1, \pm i$ ).

### 3 INVARIANT SUBMANIFOLD OF CO-DIMENSION 2 OF A MANIFOLD WITH $(\psi, g, u, v, e, \lambda)$ -STRUCTURE

We consider a submanifold  $V_m$  of  $V_n$  represented by

$$x^h = x^h(y^a) \quad \dots(3.1)$$

and put

$$B_b^h \partial_b x^h, \partial_b = \partial/\partial y^b. \quad \dots(3.2)$$

$a, b, c, \dots$  run over  $1, 2, \dots, n$  and  $h, i, j$  run over  $1$  to  $m$ .

The induced Riemannian metric is given by

$$g_{cb} = g_{it} B_c^i B_b^t. \quad \dots(3.3)$$

We denote by  $C_x^h$ ,  $2m - n$  mutually orthogonal unit normals to  $N$ . The equations of Gauss and those of Weingarten are given by (Yoshiko 1972)

$$\nabla_c B_b^h = \sum_x h_{cbx} C_x^h \quad \dots(3.4)$$

$$\nabla_c C_x^h = -h_{cx}^a B_a^h + \sum_y l_{cxy} C_x^h \quad \dots(3.5)$$

where

$$\nabla_c B_b^h = \partial_c B_b^h + \left\{ \begin{matrix} h \\ ji \end{matrix} \right\} B_c^j B_b^i - \left\{ \begin{matrix} a \\ cb \end{matrix} \right\} B_a^h \quad \dots(3.6)$$

is the Vander Warden - Borotolotti covariant differential of  $B_b^h$ ,  $\left\{ \begin{matrix} a \\ cb \end{matrix} \right\}$  being Christoffel symbols formed with  $g_{cb}$ .

$$\nabla_c C_x^h = \partial_c C_x^h + \left\{ \begin{matrix} h \\ ji \end{matrix} \right\} B_c^j C_x^i. \quad \dots(3.7)$$

$h_{cbx}$  are the components of the second fundamental tensors w.r.t. the normals  $C_x^h$ .

$$h_{cx}^a = h_{cbx} g^{ba} \quad \dots(3.8)$$

$g^{ba}$  being contravariant components of the induced metric and  $l_{cxy}$  components of the third fundamental tensor w.r.t. the normal  $C_x^h$ .

We assume that the submanifold  $V_m$  of  $V_n$  is  $\psi$ -invariant, hence we have (Yoshiko 1972, Yano and Okumura 1970)

$$\psi_i^h B_b^i = \psi_b^a B_a^h. \quad \dots(3.9)$$

$\psi_b^a$  is a (1.1) tensor field of  $V_m$ . This shows that

$$\psi_{ih} B_b^i C_x^h = 0.$$

$\psi_i^h C_x^i$  is normal to the submanifold  $V_m$ . Thus

$$\psi_i^h C_x^i = \sum_y \gamma_{xy} C_y^h.$$

Since  $\psi_{ih} C_x^i C_y^h = -\gamma_{xy}$

we see that

$$\gamma_{xy} = -\gamma_{yx}. \quad \dots(3.10)$$

Let

$$u^h = B_a^h u^a + \sum_x \alpha_x C_x^h \quad \dots(3.11)$$

$$v^h = B_a^h v^a + \sum_x \beta_x C_x^h \quad \dots(3.12)$$

and

$$e^h = B_a^h e^a + \sum_x \epsilon_x C_x^h \quad \dots(3.13)$$

$u^a, v^a, e^a$  being vector fields of  $V_m$  and  $\alpha_x, \beta_x, \epsilon_x$  are functions in  $V_m$ .

Now from (1.1), (3.9) and (3.11) to (3.12), we obtain the following relations

$$\psi_a^a \psi_b^b = -\delta_b^a + u_b u^a + v_b v^a + e_b e^a \quad \dots(3.14)$$

$$u_b \alpha_x + v_b \beta_x + e_b \epsilon_x = 0 \quad \dots(3.15)$$

$$\psi_a^a \psi_b^b g_{cd} = g_{cb} - u_b u_c - v_b v_c - e_b e_c \quad \dots(3.16)$$

$$\alpha_x u^a + \beta_x v^a + \epsilon_x e^a = 0 \quad \dots(3.17)$$

$$\sum_y \gamma_{xy} \gamma_{yz} = -\delta_{xz} + \alpha_x \alpha_z + \beta_x \beta_z + \epsilon_x \epsilon_z \quad \dots(3.18)$$

$$\psi_a^h u^a = -\lambda v^h, \psi_a^h v^a = \lambda u^h, \psi_a^h e^a = 0 \quad \dots(3.19)$$

$$\left. \begin{aligned} \sum_x \alpha_x \gamma_{xy} &= -\lambda \beta_y, \sum_x \epsilon_x \gamma_{xy} = 0 \\ \sum_x \beta_x \gamma_{xy} &= \lambda \alpha_y \end{aligned} \right\} \quad \dots(3.20)$$

$$\left. \begin{aligned} u_a u^a &= 1 - \lambda^2 - \sum_x \alpha_x^2 \\ v_a v^a &= 1 - \lambda^2 - \sum_x \beta_x^2 \\ e_a e^a &= 1 - \sum_x \epsilon_x^2 \end{aligned} \right\} \quad \dots(3.21)$$

$$\left. \begin{aligned} u_a v^a &= -\sum_x \beta_x \alpha_x \\ u_a e^a &= -\sum_x \alpha_x \epsilon_x \end{aligned} \right\} \quad \dots(3.22)$$

it can be easily seen that  $\psi_{cb}$  is skew-symmetric, since

$$\psi_{ji} B_o^j B_o^i = \psi_{cb}. \quad \dots(3.23)$$

Equations (3.14) – (3.22) show that a necessary and sufficient condition for  $\psi_b^a, g_{cb}, u_b, v_b, e_b$  and  $\lambda$  to admit a  $(\psi, g, u, v, e, \lambda)$ -structure is that  $\sum_x \alpha_x^2 = 0, \sum_x \beta_x^2 = 0$  and  $\sum_x \epsilon_x^2 = 0$ .  $\alpha_x, \beta_x$  and  $\epsilon_x$  are zero or we say that  $u^h, v^h$  and  $e^h$  are always tangent to the submanifold. Since

$$(\nabla_i u_j - \nabla_j u_i) B_c^j B_b^i = \nabla_c u_b - \nabla_b u_c \tag{3.24a}$$

and

$$(\nabla_j u_i + \nabla_i u_j) B_c^j B_b^i = \nabla_c u_b + \nabla_b u_c \tag{3.24b}$$

hence

$$\begin{aligned} S_{ji}^h B_c^j B_b^i &= (N_{cb}^a + u_{cb}u^a + v_{cb}v^a + e_{cb}e^a) B_a^h \\ &+ \left\{ \sum_x (\alpha_x u_{cb} + v_{cb} \beta_x + e_{cb} \epsilon_x) \right\} C_x^h. \end{aligned} \tag{3.25}$$

where  $N_{cb}^a$  is the Nijenhuis tensor  $\psi_b^a$ .

This leads to the fact that if manifold with  $(\psi, g, u, v, e, \lambda)$ -structure is normal then the induced structure  $(\psi, g, u, v, e, \lambda)$  on the submanifold  $V_m$  is also normal.

In view of (3.24a, b) we can also say that  $u^a, v^a, e^a$  will be harmonic vector fields on induced submanifold, if  $u^h, v^h$  and  $e^h$  are the harmonic vectors on the manifold  $V_n$  with  $(\psi, g, u, v, e, \lambda)$ -structure. Therefore

*Theorem 3.1* — The induced  $(\psi, g, u, v, e, \lambda)$ -structure of the invariant submanifold  $V$  of co-dimension 2 with quasi-normal  $(\psi, g, u, v, e, \lambda)$ -structure is also quasi-normal.

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