

# SOME THEOREMS ON AFFINE MOTION IN A RECURRENT FINSLER SPACE

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In the present paper an infinitesimal affine motion in a recurrent Finsler space is considered. In such space many important results have been obtained.

## 1. INTRODUCTION

Let us consider an  $n$ -dimensional Finsler space  $F_n$  (Rund 1959) having positively homogeneous metric function  $F(x, \dot{x})$  of degree one in  $\dot{x}^i, s$ . The fundamental metric tensor of the space is defined by

$$g_{ij}(x, \dot{x}) = \frac{1}{2} \dot{\mathfrak{D}}_j \dot{\mathfrak{D}}_i F^2(x, \dot{x}), \quad (\dot{\mathfrak{D}}_i \equiv \mathfrak{D} / \mathfrak{D} \dot{x}^i) \quad \dots(1.1)$$

which is symmetric in its lower indices.

Let  $X^i(x, \dot{x})$  be any contravariant vector field depending both upon positional and directional arguments. Cartan's covariant derivative of  $X^i(x, \dot{x})$  with respect to  $x^k$  is given by (Rund 1959)

$$X^i{}_{;k} = \mathfrak{D}_k X^i - (\dot{\mathfrak{D}}_m X^i) G^m{}_k + X^h \Gamma^*{}_{hk}{}^i \quad \dots(1.2)$$

where  $\Gamma^*{}_{hk}{}^i(x, \dot{x})$  are connection coefficients.

The curvature tensor  $K^i{}_{hjk}(x, \dot{x})$  arising from this covariant is given by

$$K^i{}_{hjk}(x, \dot{x}) = 2 \left\{ \mathfrak{D}_{[k} \Gamma^*{}_{j]h}{}^i - \dot{\mathfrak{D}}_r \Gamma^*{}_{hlj}{}^i G^r{}_{kl} + \Gamma^*{}_{hlj}{}^r \Gamma^*{}_{klr}{}^i \right\} \quad \dots(1.3)$$

Let us consider an infinitesimal point transformation

$$\bar{x}^i = x^i + v^i(x) dt \quad \dots(1.4)$$

where  $v^i(x)$  is any vector field and  $dt$  an infinitesimal constant. In view of the covariant derivative and (1.4) the Lie-derivatives of any tensor field

$$2A_{[hk]} = A_{hk} - A_{kh} \text{ and } 2A_{(hk)} = A_{hk} + A_{kh}.$$

$T_j(x, \dot{x})$  and connection coefficient  $\Gamma_{jk}^{*i}(x, \dot{x})$  are given by (Yano 1957)

$$Lv T_j^i(x, \dot{x}) = T_{j|h}^i v^h + (\dot{\nabla}_s T_j^i) v_{|r}^s \dot{x}^r + T_h^i v_{|j}^h - T_j^h v_{|h}^i \quad \dots(1.5)$$

$$Lv \Gamma_{jk}^{*i}(x, \dot{x}) = v_{|ijk}^i + K_{jkh}^i v^h + (\dot{\nabla}_s \Gamma_{jk}^{*i}) v_{|r}^s \dot{x}^r \quad \dots(1.6)$$

We have the following commutation formulae :

$$\dot{\nabla}_l (Lv T_j^i) - Lv \dot{\nabla}_l T_j^i = 0 \quad \text{and} \quad \dots(1.7)$$

$$(Lv \Gamma_{jh}^{*i})_{|k} - (Lv \Gamma_{kh}^{*i})_{|j} = Lv K_{hjk}^i + 2 \dot{x}^s \dot{\nabla}_r \Gamma_{h|j}^{*i} Lv \Gamma_{k|s}^{*r} \quad \dots(1.8)$$

If in a Finsler space  $F_n$  the curvature tensor  $K_{hjk}^i(x, \dot{x})$  satisfies the relation

$$K_{hjk|lm}^i = \beta_m K_{hjk}^i \quad \dots(1.9)$$

where  $\beta_m(x)$  is any non-zero vector then such type of  $F_n$  is called recurrent Finsler space. We shall denote it by  $F_n^*$  throughout this paper.

### 2. AFFINE MOTION IN $F_n^*$

With the help of the infinitesimal point transformation (1.4) we can obtain a deformed space with affine connection  $\Gamma_{jk}^{*i} + (Lv \Gamma_{jk}^{*i}) dt$ .

If the original space and deformed space have the same affine connection, the transformation is called affine motion of the space  $F_n$ . In our  $F_n^*$  space we shall consider such type of affine motion. In order that it be the case, it is necessary and sufficient that we have

$$Lv \Gamma_{jk}^{*i} = v_{|ijk}^i + K_{jkh}^i v^h + (\dot{\nabla}_s \Gamma_{jk}^{*i}) v_{|r}^s \dot{x}^r = 0 \quad \dots(2.1)$$

In view of (2.1), eqn. (1.8) reduces to

$$Lv K_{jkh}^i = 0. \quad \dots(2.2)$$

With respect to the Lie-derivative, for any tensor  $T_{jk}^i$ , we have

$$\begin{aligned} (Lv T_{jk|l}^i) - (Lv T_{j|kl}^i) &= T_{sk}^i Lv \Gamma^{*s}_{jl} + T_{j|l}^i Lv \Gamma^{*s}_{kl} - T_{j|k}^s Lv \Gamma^{*i}_{sl} \\ &+ (\dot{\nabla}_s T_{jk}^i) Lv \Gamma^{*s}_{rl} \dot{x}^r \end{aligned} \quad \dots(2.3)$$

Apply this commutation formula to the curvature tensor  $K_{hjk}^i(x, \dot{x})$  and noting the equations (2.1) and (2.2), we get



$\alpha(x)$  has to vanish, say  $v^i$  has to span a field of parallel contravariant vectors (a contrafield).

4. GENERAL CASE

In view of (1.5), the Lie-derivative of the curvature tensor field  $K^i_{hjk}(x, \dot{x})$  is given by

$$Lv K^i_{hjk}(x, \dot{x}) = K^i_{hjkls} v^s + (\dot{\nabla}_s K^i_{hjk}) v^s_n x^r + K^i_{sjk} v^s_{|h} + K^i_{hsk} v^s_{|j} + K^i_{hjs} v^s_{|k} - K^s_{hjk} v^i_{|s} \dots(4.1)$$

Introducing (1.9) and (2.2) into the above equation, we get

$$\beta_s v^s K^i_{hjk} + (\dot{\nabla}_s K^i_{hjk}) v^s_{|r} x^r + K^i_{sjk} v^s_{|h} + K^i_{hsk} v^s_{|j} + K^i_{hjs} v^s_{|h} - K^s_{hjk} v^i_{|s} = 0 \dots(4.2)$$

In  $Fn^*$ , let us consider an affine motion of the following form

$$\overline{x^j} = x^j + v^j(x) dt, \quad v^i_{|j} = c \delta^i_j \dots(4.3)$$

where  $c(x) \neq 0$ .

In view of (4.3), equ. (4.2) yields

$$(v^s \beta_s + 2c) K^i_{hjk} = 0 \dots(4.4)$$

Since  $K^i_{hjk} \neq 0$ , from this, we get

$$c = -\frac{1}{2} v^s \beta_s \dots(4.5)$$

Hence the motion (4.3) takes the following form

$$\overline{x^i} = x^i + v^i(x) dt, \quad v^i_{|j} = -\frac{1}{2} v^s \beta_s \delta^i_j \dots(4.6)$$

Substituting the latter part of (4.6) into (2.1), we obtain

$$K^i_{jkh} v^h = \frac{1}{2} \{ 2\beta_{m|k} - \beta_m \beta_k \} v^m \delta^i_j \dots(4.7)$$

We have the following identities (Rund 1959):

$$K^i_{hjk|lm} + K^i_{hkm|lj} + K^i_{hmi|jk} = -x^s \{ (\dot{\nabla}_l \Gamma^{*i}_{hj}) K^l_{skm} + (\dot{\nabla}_l \Gamma^{*i}_{hm}) K^l_{sjk} + (\dot{\nabla}_l \Gamma^{*i}_{hk}) K^l_{smj} \} \dots(4.8)$$

and

$$K^i_{hjk} + K^l_{jkh} + K^l_{khj} = 0 \dots(4.9)$$

Introducing (1.9) into (4.8), we get

$$\beta_m K^l_{hjk} + \beta_j K^l_{hkm} + \beta_k K^l_{hmi} = -x^s \{ (\dot{\nabla}_l \Gamma^{*i}_{hj}) K^l_{skm} + (\dot{\nabla}_l \Gamma^{*i}_{hm}) K^l_{sjk} + (\dot{\nabla}_l \Gamma^{*i}_{hk}) K^l_{smj} \} \dots(4.10)$$

Multiplying this identity by  $v^m$  and summing up with respect to  $m$ , we obtain

$$K^i_{hjk} \beta_m v^m = K^i_{hjm} v^m \beta_k - \beta_j K^i_{hkm} v^m - \dot{x}^s \{ (\dot{\mathfrak{D}}_l \Gamma^{*i}_{hj}) K^l_{skm} v^m + (\dot{\mathfrak{D}}_l \Gamma^{*i}_{hm}) K^l_{sjk} v^m - (\dot{\mathfrak{D}}_m \Gamma^{*i}_{hk}) K^l_{sjm} v^m \} \tag{4.11}$$

where we have used  $K^i_{hjk} = -K^i_{hkj}$

Introducing (4.7) into the right-hand side of (4.11), we get

$$K^i_{hjk} \beta_m v^m = \frac{1}{2} \{ \beta_{rj} \beta_k - \beta_{r,k} \beta_j \} v \delta^i_h - \dot{x}^s (\dot{\mathfrak{D}}_l \Gamma^{*i}_{hm}) K^l_{sjk} v^m \tag{4.12}$$

Multiplying the identity (4.9) by  $\beta_m v^m$ , we have

$$K^i_{hjk} \beta_m v^m + K^i_{jkh} \beta_m v^m + K^i_{khj} \beta_m v^m = 0 \tag{4.13}$$

Substituting (4.12) into the above identity, we obtain

$$\frac{1}{2} v^r [ \delta^i_h (\beta_k \beta_{rj} - \beta_{r,k} \beta_j) + \delta^i_j (\beta_{r,k} \beta_h - \beta_{r,h} \beta_k) + \delta^i_k (\beta_{r,h} \beta_j - \beta_{r,j} \beta_h) ] - \dot{x}^s v^m \{ (\dot{\mathfrak{D}}_l \Gamma^{*i}_{hm}) K^l_{sjk} + (\dot{\mathfrak{D}}_l \Gamma^{*i}_{jm}) K^l_{skh} (\dot{\mathfrak{D}}_l \Gamma^{*i}_{km}) K^l_{shj} \} = 0 \tag{4.14}$$

Contracting the above equation with respect to the indices  $i$  and  $h$ , we get

$$\frac{(n-2)}{2} (\beta_{rj} \beta_k - \beta_{r,k} \beta_j) v^r - \dot{x}^s v^m \{ (\dot{\mathfrak{D}}_l \Gamma^{*i}_{im}) K^l_{sjk} + (\dot{\mathfrak{D}}_l \Gamma^{*i}_{jm}) K^l_{skj} + (\dot{\mathfrak{D}}_l \Gamma^{*i}_{km}) K^l_{sij} \} = 0 \tag{4.15}$$

Comparing this result with (4.12) and after little simplification we obtain

$$(n-2) K^i_{hjk} \beta_m v^m = \dot{x}^s v^m \{ (n-1) (\dot{\mathfrak{D}}_l \Gamma^{*i}_{hm}) K^l_{sjk} + (\dot{\mathfrak{D}}_l \Gamma^{*i}_{jm}) K^l_{skh} + (\dot{\mathfrak{D}}_l \Gamma^{*i}_{km}) K^l_{shj} \} \tag{4.16}$$

Thus, we have the following theorem.

*Theorem 4.1* — If an  $F_n^*$  admits an affine motion of the form (4.3) then the equation (4.16) holds.

In a general affinely connected Finsler space  $F_n$ , if  $v^i_{|i} = 0$ , the vector  $v^i$  determines a contrafield. We shall also consider a contrafield spanned by  $v^i(x)$  in our  $F_n^*$ .

In  $F_n^*$ , we shall consider to find a necessary and sufficient condition upon the possibility of a special affine motion of the form

$$\bar{x}^i = x^i + v^i(x) dt v^i_{|j} = 0 \tag{4.17}$$

Now, if it will be the case, the eqn. (2.1) of motion becomes

$$L v \Gamma^{*i}_{jk} = K^i_{jkh} v^h = 0 \tag{4.18}$$

and the integrability condition of (2.1) becomes

$$L v K^i_{hjk} = K^i_{hjk} \beta_s v^s = 0 \tag{4.19}$$

So, owing to the non-flatness of the space, from (4.19), we get

$$\beta_s v^s = 0 \quad \dots(4.20)$$

In view of (4.18) and (4.20), we can conclude as follows :

If an  $F_n^*$  space admits an affine motion of the form (4.17), it is necessary that we have

$$v^s \beta_s = 0, K^i_{jkh} v^h = 0. \quad \dots(4.21)$$

Conversely if we have (4.21), using the formula (4.9) and the latter part of (4.21), we get

$$K^i_{hjk} v^h = 0$$

From the general theory of parallel vectors (Knebelman 1945),  $v^i$  determines a contrafield in  $F_n^*$ ,  $v^i_{;j} = 0$ , when and only when  $v^i$  satisfy the equation (4.21). Thus, we may regard  $v^i(x)$  as contrafield in  $F_n^*$ . In such a case  $LvK^i_{hjk} = 0$  holds good and which is the integrability condition of  $Lv\Gamma^{*i}_{jk} = 0$ . Because, if  $v^i_{;j} = 0$ , by virtue of its definition,  $LvK^i_{hjk}$  is written as  $v^s \beta_s K^i_{hjk}$ . However, we have the former of (4.21) so we get  $LvK^i_{hjk} = 0$ , consequently  $Lv\Gamma^{*i}_{jk} = 0$  are integrable. Thus, (4.21) is also a sufficient condition for  $F_n^*$  admitting (4.17).

Thus, we have the following theorem.

*Theorem 4.2* — The necessary and sufficient condition that a non-flat  $F_n^*$  admits an infinitesimal transformation  $\bar{x}^i = x^i + v^i(x) dt$  in order that the vector  $v^i(x)$  span a contrafield in  $F_n^*$  is that eqn. (4.21) be valid.

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