

# COSMOLOGICAL SOLUTION OF EINSTEIN'S EQUATIONS IN GENERAL RELATIVITY

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A model is proposed in which the central pressure and central density are finite.

## 1. INTRODUCTION

Krori *et al.* (1974) found an internal solution for a spherically symmetric matter distribution taking

$$\rho c^2 + 3p = f(r).$$

Here in the present paper, we have attempted a solution of the Einstein's field equations taking cosmological constant in the solution. The constant  $k_1$  is any arbitrary positive value of the order of  $\Lambda$ ,  $k_3$  is negative and  $k_2$  is positive lying between 0 and 1. With positive value of  $\Lambda$ , it is found that radius increases and consequently density at the centre increases to counteract this expansion.

## 2. FIELD EQUATIONS AND THEIR SOLUTIONS

The line element exhibiting spherically symmetric matter distribution (Møller 1952, p. 329),

$$ds^2 = a(r) dr^2 + r^2 (d\theta^2 + \sin^2 \theta d\phi^2) - b(r) c^2 dt^2. \quad \dots(2.1)$$

The equations to be satisfied are

$$p' + (\rho c^2 + p) \frac{b'}{2b} = 0 \quad \dots(2.2)$$

$$\frac{b'}{abr} - \frac{1}{r^2} \left(1 - \frac{1}{a}\right) + \Lambda = 8\pi p \quad \dots(2.3)$$

$$\frac{a'}{a^2 r} + \frac{1}{r^2} \left(1 - \frac{1}{a}\right) - \Lambda = 8\pi \rho c^2 \quad \dots(2.4)$$

where prime denotes differentiation with reference to  $r$  and  $\Lambda$  is the cosmological constant. Now, we assume the solution

$$b = \frac{1}{2} k_1 r^2 + k_2 \quad \dots(2.5)$$

where  $k_1$  and  $k_2$  are constants of integration. Following Krori *et al.* (1974), we have,

$$p = \frac{1}{16\pi} \left( \frac{b'}{abr} \right) + k_3 \quad \dots(2.6)$$

where  $k_3$  is a constant of integration.

From equation (2.3)

$$a = \frac{k_1 r^2 + k_2}{(1 + 8\pi k_3 r^2 - r^2 \Lambda) (\frac{1}{2} k_1 r^2 + k_2)} \quad \dots(2.7)$$

Hence from equation (2.6)

$$p = \frac{1}{16\pi} \left( \frac{1 + 8\pi k_3 r^2 - r^2 \Lambda}{r^2 + \frac{k_2}{k_1}} \right) + k_3 \quad \dots(2.8)$$

and from (2.4)

$$8\pi\rho c^2 = \frac{\frac{k_1^2}{2} (\Lambda - 24\pi k_3) r^4 + \frac{1}{2} k_1 r^2 (k_1 + 3k_2 \Lambda - 56\pi k_2 k_3) + k_2 (\frac{3}{2} k_1 + 2k_2 \Lambda - 24\pi k_2 k_3)}{(k_1 r^2 + k_2)^2} \quad \dots(2.9)$$

The boundary is given by

$$r_1 = \left\{ \frac{k_1 + 16\pi k_2 k_3}{k_1 (\Lambda - 24\pi k_3)} \right\}^{1/2} \quad \dots(2.10)$$

At the boundary, since  $a(r_1) b(r_1) = 1$ , we have

$$(k_1 + \Lambda + 8\pi k_3) r_1^2 = (1 - k_2). \quad \dots(2.11)$$

Putting the value of  $r_1$  in equation (2.11) and simplifying, we have

$$k_1 = -16\pi k_3. \quad \dots(2.12)$$

Thus if  $k_1$  is positive and of the order of  $\Lambda$ ,  $k_3$  will be negative, hence from (2.11)  $k_2$  lies between 0 and 1.

Now,

$$b(r_1) = 1 - \frac{2m}{r_1} - \frac{1}{3} \Lambda r_1^2$$

at the exterior gives

$$\frac{2m}{r_1} = \frac{1}{2} k_1 r_1^2 + \frac{2}{3} r_1^2 (\Lambda - 12\pi k_3). \quad \dots(2.13)$$

Since  $k_1$  is any arbitrary positive value of the order of  $\Lambda$  and since  $k_3$  is negative,  $2m/r_1$  is positive.

Equation (2.13) gives using the value of  $k_3$  from equation (2.12)

$$k_1 = \frac{2m}{r_1^3} - \frac{2}{3}\Lambda \quad \dots(2.14)$$

where  $\Lambda$  (lower limit)  $\simeq 9.3 \times 10^{-58} \text{ cm}^{-2}$  (Tolman 1964, p. 345). Since  $k_1$  is considered a positive value, hence  $m/r_1^3$  must be greater than  $\frac{1}{3}\Lambda$ .

At  $r = 0$ , we have

$$8\pi\rho_0 c^2 = \frac{3k_1}{2k_2} + 2(\Lambda - 12\pi k_3) \quad \dots(2.15)$$

$$8\pi p_0 = \frac{k_1}{2k_2} + 8\pi k_3. \quad \dots(2.16)$$

From above it is found that  $\rho_0 c^2$  and  $p_0$  are positive and  $\rho_0 c^2 > 3p_0$ . Also we have

$$\frac{\rho_0 c^2}{p_0} = 3 \left( \frac{1+k_2}{1-k_2} \right) + \frac{4k_2\Lambda}{\left( \frac{2m}{r_1^3} - \frac{2}{3}\Lambda \right) (1-k_2)}. \quad \dots(2.17)$$

The equation (2.10) can be simplified to as

$$r_i = \left( \frac{1-k_2}{\Lambda + \frac{8}{3}k_1} \right)^{1/2} \quad \dots(2.18)$$

and it is real since  $k_1$  is positive and  $k_2$  lies between 0 to 1.

Since  $\Lambda$  is taken as a positive value, the mass density at the centre is slightly increased [eqn. (2.15)] and this is necessary to counteract any expansion [eqn. (2.18)] that arises due to the inclusion of cosmological term  $\Lambda$  in the solution.

#### REFERENCES

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