

ON SIMILARITY SOLUTION OF CHEMICAL BOUNDARY LAYER THEORY

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This paper presents an investigation of similarity solutions in chemical boundary layer theory. It is seen that when the flow is accompanied with n th order homogeneous volume reaction, we may get similarity solution when $n \neq 1$.

INTRODUCTION

There is a possibility of an interesting and important investigation in boundary layer theory when the flow is along a soluble body and accompanied by chemical reactions. This type of study is useful in chemical physics and engineering. Some research work in this direction has been done by Ghosal *et al.* (1972) and others.

ANALYSIS

Let us consider the flow along a flat plate that contains a species A slightly soluble in the fluid B . The concentration at the plate surface is assumed to be \bar{C}_{Ae} which is the solubility of A in B and the concentration far from the plate is assumed to be $\bar{C}_{A\infty} = 0$.

We are considering here the case when A reacts with B by an n th order homogeneous chemical reaction (volume). We further assume here that the concentration of dissolved A is small enough to consider the physical properties ρ , μ , D_{AB} to be constants throughout the fluid, where ρ and μ are respectively the density and the viscosity of the fluid B and D_{AB} is the diffusivity of A in B .

The boundary layer equations of the above system are adequately described by the following equations (Bird *et al.* 1960):

$$\frac{\partial \bar{u}}{\partial \bar{t}} + \bar{u} \frac{\partial \bar{u}}{\partial \bar{x}} + \bar{v} \frac{\partial \bar{u}}{\partial \bar{y}} = \nu \frac{\partial^2 \bar{u}}{\partial \bar{y}^2} \quad \dots(1)$$

$$\frac{\partial \bar{u}}{\partial \bar{x}} + \frac{\partial \bar{v}}{\partial \bar{y}} = 0 \quad \dots(2)$$

$$\frac{\partial \bar{C}}{\partial \bar{t}} + \bar{u} \frac{\partial \bar{C}}{\partial \bar{x}} + \bar{v} \frac{\partial \bar{C}}{\partial \bar{y}} = D_{AB} \frac{\partial^2 \bar{C}}{\partial \bar{y}^2} + K \bar{C}^n \quad \dots(3)$$

where \bar{C} is the concentration of the solute which is a function of $\bar{x}, \bar{y}, \bar{t}$; n is an integer which indicates the order of the reaction and K the volume rate constant.

The boundary conditions of the system are given by

$$\left. \begin{aligned} \bar{y} = 0; \quad \bar{u} = \bar{v} = 0, \quad \bar{C} = \bar{C}_w(\bar{x}, \bar{t}) \\ \bar{y} \rightarrow \infty; \quad \bar{u} \rightarrow \bar{U}_\infty(\text{constant}), \quad \bar{C} \rightarrow 0 \end{aligned} \right\} \quad \dots(4)$$

(subscript w is the wall condition).

The variables u, v, x, y, t are made dimensionless as follows:

$$\left. \begin{aligned} x = \frac{\bar{x}}{L}, \quad y = \frac{\bar{y}}{L}, \quad t = \frac{\bar{t}}{L^2} \\ u = \frac{\bar{u}L}{v}, \quad v = \frac{\bar{v}L}{v}, \quad C = \frac{\bar{C}}{C_0} \end{aligned} \right\} \quad \dots(5)$$

where L is a representative length and C_0 denotes a typical value of \bar{C} , say at $x = 0$ on the plate at $t = 0$. Introducing (5) in eqns. (1), (2) and (3) we get

$$\frac{\partial u}{\partial t} + u \frac{\partial u}{\partial x} + v \frac{\partial u}{\partial y} = \frac{\partial^2 u}{\partial y^2} \quad \dots(6)$$

$$\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} = 0 \quad \dots(7)$$

$$\frac{\partial C}{\partial t} + u \frac{\partial C}{\partial x} + v \frac{\partial C}{\partial y} = \frac{1}{Sc} \frac{\partial^2 C}{\partial y^2} + \frac{C_0^{n-1} L^2 K}{v} C^n \quad \dots(8)$$

where

$$Sc = \frac{v}{D_{AB}}$$

is the Schmidt's number and

$$\frac{C_0^{n-1} L^2 K}{v}$$

is a dimensionless constant.

The boundary conditions now become

$$\left. \begin{aligned} y = 0; \quad u = v = 0, \quad C = C_w(x, t) \\ y \rightarrow \infty; \quad u \rightarrow U_\infty(\text{constant}), \quad C \rightarrow 0 \end{aligned} \right\} \quad \dots(9)$$

The equation of continuity is satisfied if we choose a dimensionless stream function ψ such that

$$u = \frac{\partial \psi}{\partial y}, \quad v = -\frac{\partial \psi}{\partial x}.$$

If we choose, a similarity variable in the form

$$\eta = y \phi_1(x, t) \quad \dots(10)$$

and the dependent variables as follows:

$$f(\eta) = \psi(x, y, t)/\phi_2(x, t) \quad \dots(11)$$

$$\theta(\eta) = AC(x, y, t)/C_w(x, t) \quad \dots(12)$$

(where A is a constant which is to be determined), the present problem reduces to one of determining the functions $\phi_1(x, t)$, $\phi_2(x, t)$ and $C_w(x, t)$ such that eqns. (6) and (8) are reducible to ordinary differential equations for $f(\eta)$ and $\theta(\eta)$ with proper boundary conditions. With the following identities from eqns. (10), (11) and (12),

$$u = \frac{\partial \psi}{\partial y} = \phi_1 \phi_2 f'(\eta) \quad \dots(13)$$

$$v = -\frac{\partial \psi}{\partial x} = -\left[\frac{\partial \phi_2}{\partial x} f + \frac{\phi_2}{\phi_1} \frac{\partial \phi_1}{\partial x} \eta f' \right] \quad \dots(14)$$

$$C = \frac{C_w}{A} \theta \quad \dots(15)$$

substituted in eqns. (6) and (8) we obtain

$$f'''' - a_2 \eta f'' - (a_2 + a_3) f' + a_5 f f'' - (a_4 + a_5) f'^2 = 0 \quad \dots(16)$$

$$\frac{1}{Sc} \theta'' - (a_2 \eta - a_5 f) \theta' - (a_6 + a_7 f') \theta + a_1 \theta^n = 0 \quad \dots(17)$$

where

$$a_1 = C_w^{n-1} / \phi_1^2 \quad \dots(18)$$

$$a_2 = \phi_{1x} / \phi_1^3 \quad \dots(19)$$

$$a_3 = \phi_{2t} / (\phi_1^2 / \phi_2) \quad \dots(20)$$

$$a_4 = (\phi_2 \phi_{1x}) / \phi_1^2 \quad \dots(21)$$

$$a_5 = \phi_{2x} / \phi_1 \quad \dots(22)$$

$$a_6 = C_{wt} / (C_w \phi_1^2) \quad \dots(23)$$

$$a_7 = \phi_2 C_{wx} / (\phi_1 C_w) \quad \dots(24)$$

It is to be noted here that A is made equal to $C_0 (L^2 k/\nu)^{1/(n-1)}$. We now try to find out for what values of n similarity solutions for eqns. (16) and (17) are possible. From eqn. (19) we get

$$\phi_1 = [B(x) - 2a_2 t]^{-\frac{1}{2}} \quad \dots(25)$$

where $B(x)$ is yet to be determined. Combining eqns. (19) and (20) we get

$$\phi_2 = D(x) \phi_1^n \quad \dots(26)$$

where $D(x)$ is an arbitrary function and $w = a_3/a_2$.

Thus

$$\phi_2 = D(x) [B(x) - 2a_2t]^{-w/2}. \quad \dots(27)$$

A similar combination of eqns. (21) and (22) gives

$$\phi_2 = E(t) \phi_1^\epsilon \quad \dots(28)$$

where $\epsilon = a_5/a_4$, $a_4 \neq 0$. The case of $a_4 = 0$ will be treated later. Equation (21) may be solved giving

$$\phi_1 = \left[\frac{a_4(\epsilon - 1)}{E(t)} x + F(t) \right]^{1/(\epsilon-1)} \quad \dots(29)$$

and therefore ϕ_2 is given by

$$\phi_2 = E(t) \left[\frac{a_4(\epsilon - 1)}{E(t)} x + F(t) \right]^{\epsilon/(\epsilon-1)} \quad \dots(30)$$

Comparing eqns. (25), (27), (29) and (30), we get

$$\epsilon = w = -1, \quad B(x) = -\frac{2a_4}{a_{12}} x$$

$$D(x) = E(t) = a_{12}, \quad F(t) = -2a_2t.$$

With the above results we get

$$\phi_1 = \left(-\frac{2a_4}{a_{12}} x - 2a_2t \right)^{-1/2} \quad \dots(31)$$

$$\phi_2 = a_{12} \left(-\frac{2a_4}{a_{12}} x - 2a_2t \right)^{1/2} \quad \dots(32)$$

From relations (18) and (31) it appears that

$$\begin{aligned} C_w &= a_1^{1/(n-1)} \left(-\frac{2a_4}{a_{12}} x - 2a_2t \right)^{-1/(n-1)} \\ &= a_{13} \left(-\frac{2a_4}{a_{12}} x - 2a_2t \right)^{-1/(n-1)} \end{aligned} \quad \dots(33)$$

where $a_{13} = a_1^{1/(n-1)}$.

Now it is seen that the right-hand sides of relations (18)–(24) become constant with the above results for ϕ_1 , ϕ_2 and C_w . We see that n may be any value except 1. Thus similarity solutions of eqns. (16) and (17) are possible if $n \neq 1$.

The case of $a_4 = 0$ will be considered now. If $a_4 = 0$ then eqn. (25) becomes

$$\phi_1 = [a_{14} - 2a_2t]^{-\frac{1}{2}} \quad \dots(34)$$

where a_{14} is a constant of integration.

From eqns. (20), (22) and (34) we ultimately get

$$\phi_2 = (a_5x + a_{15})(a_{14} - 2a_2t)^{-\frac{1}{2}} \quad \dots(35)$$

where a_{15} is a constant.

Also from eqns. (18) and (34) we get

$$C_w = b(a_{14} - 2a_2t)^{-1/(n-1)}$$

where $b = a_1^{1/(n-1)}$.

It can be easily verified that the right-hand sides of relations (23) and (24) are constants with the above results for ϕ_1 , ϕ_2 and C_w . Hence similarity solution is also possible in this case for any value of n except $n = 1$.

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