

# ON THE NON-SUMMABILITY OF THE CONJUGATE SERIES OF A FOURIER SERIES

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1. Let  $f(\theta)$  be a periodic function with period  $2\pi$  and integrable ( $L$ ) in  $(-\pi, \pi)$ . Let the Fourier series of  $f(\theta)$  be

$$(1.1) \quad \frac{1}{2}a_0 + \sum_{n=1}^{\infty} (a_n \cos n\theta + b_n \sin n\theta).$$

Then the conjugate series of the Fourier series is given by

$$(1.2) \quad \sum_{n=1}^{\infty} B_n(\theta) = \sum_{n=1}^{\infty} (b_n \cos n\theta - a_n \sin n\theta),$$

and the conjugate function, defined as a Cauchy-integral, is given by

$$(1.3) \quad g(\theta) = (1/2\pi) \lim_{\epsilon \rightarrow 0} \int_{\epsilon}^{\pi} \psi(u) \cot (u/2) du,$$

where

$$\psi(u) = f(\theta + u) - f(\theta - u).$$

The Abel-limit associated with the series (1.2) is given by

$$(1.4) \quad \lim_{x \rightarrow 1-0} \sum_{n=1}^{\infty} (b_n \cos n\theta - a_n \sin n\theta) x^n.$$

2. In what follows we use the following notations—

$$(2.1) \quad \left\{ \begin{array}{l} \Psi_{\alpha}(t) = \frac{1}{\Gamma(\alpha)} \int_0^t (t-u)^{\alpha-1} \psi(u) du, \alpha > 0; \\ \Psi_0(t) = \psi(t); \\ \psi_{\alpha}(t) = \Gamma(\alpha+1) t^{-\alpha} \Psi_{\alpha}(t), \alpha \geq 0. \end{array} \right.$$

$$(2.2) \quad \left\{ \begin{array}{l} \Delta = 1 - 2x \cos t + x^2, \\ Q'(t) = (x \sin t) / \Delta, \\ Q(t) = t / (p^2 + t^2), \text{ where } p = \log x, \\ q(t) = (1/t) - Q(t). \end{array} \right.$$

$$(2.3) \quad \left\{ \begin{array}{l} Q^{\rho}(t) = (d/dt)^{\rho} \{ Q(t) \}, \\ q^{\rho}(t) = (d/dt)^{\rho} \{ q(t) \}. \end{array} \right.$$

$$(2.4) \quad \epsilon = \arcsin(1-x).$$

3. The following theorem on the summability ( $A$ ) of a conjugate series for which the conjugate function exists as a Cauchy-integral is due to Plessner [1923]:

*Theorem A*: Let

$$V(x, \theta) = \sum_{n=1}^{\infty} (b_n \cos n\theta - a_n \sin n\theta) x^n, \quad (0 \leq x < 1).$$

Then if for any  $\theta$ , the condition

$$(3.1) \quad \int_0^t \psi(t) dt = o(t), \quad (t \rightarrow 0)$$

is satisfied, then

$$\lim_{x \rightarrow 1} \left[ V(x, \theta) - (1/2\pi) \int_{\epsilon}^{\pi} \psi(t) \cot(t/2) dt \right] = 0.$$

It follows immediately from this theorem that if condition (3.1) holds, then the divergence to  $+\infty(-\infty)$  of the conjugate function (1.3) is a necessary and sufficient condition for the divergence of the Abel-limit of the conjugate series to  $+\infty(-\infty)$ .

Prasad [1932] obtained a more general theorem which included Plessner's theorem as a particular case. In Prasad's theorem the conjugate function was replaced by his 'Generalized Conjugate Function' and the condition (3.1) was replaced by the less stringent condition, viz.

$$\int_0^t \{\Psi(t)/t\} dt = o(t), \quad \text{as } t \rightarrow 0.$$

Generalizing Plessner's theorem in another direction, viz. by replacing the condition (3.1) by the less stringent condition

$$(3.2) \quad \psi_{\alpha}(t) = o(1), \quad \text{as } t \rightarrow 0,$$

$\alpha$  being any positive integer  $\geq 1$ , the author had proved the following [Sinha, 1953]:

*Theorem B*: If  $\psi_{\alpha}(t) = O(1)$ , as  $t \rightarrow 0$ ,  $\alpha$  being a positive integer  $\geq 1$ , then the divergence of the generalized conjugate function

$$(1/2\pi) \int_{\epsilon}^{\pi} \psi_{\alpha-1}(t) \cot(t/2) dt$$

to  $+\infty(-\infty)$  is a necessary and sufficient condition for the divergence to  $+\infty(-\infty)$  of the Abel-limit of the conjugate series of a Fourier series.

4. The object of the present paper is to extend the result of Theorem B and establish the following theorem for the more abstruse case in which  $\alpha (> 0)$  is not necessarily an integer, but may become non-integral.

*Theorem 1*: If  $\psi_{\alpha}(t) = o(1)$ , as  $t \rightarrow 0$ , where  $\alpha \geq 1$ ,

$$(4.1) \quad \left[ V(x, \theta) - (1/2\pi) \int_{\epsilon}^{\pi} \psi_{\alpha-1}(t) \cot(t/2) dt \right] = o(1), \quad \text{as } x \rightarrow 1-0.$$

*Theorem 2*: If  $\psi_{\alpha}(t) = O(1)$ , as  $t \rightarrow 0$ , where  $\alpha > 0$ ,

$$(4.2) \quad \left[ V(x, \theta) - (1/2\pi) \int_{\epsilon}^{\pi} \psi_{\alpha-1}(t) \cot(t/2) dt \right] = O(1), \quad \text{as } x \rightarrow 1-0.$$

Obviously Theorem 2 is equivalent to the statement that if  $\psi_\alpha(t) \ll O(1)$ , as  $t \rightarrow 0$ , where  $\alpha > 0$ , then the divergence of the generalized conjugate function

$$(1/2\pi) \int_{\epsilon}^{\pi} \psi_{\alpha-1}(t) \cot(t/2) dt$$

to  $+\infty$  ( $-\infty$ ) is a necessary and sufficient condition for the divergence to  $+\infty$  ( $-\infty$ ) of the Abel-limit of the conjugate series of a Fourier series.

5. For proving the above theorems, we shall repeatedly use a number of results, which may be stated below as lemmas, the proofs of most of them being quite straightforward:—

*Lemma 1 :*

$$\begin{aligned} Q^{(r)}(t) &= O(1/t^{r+1}), \quad (t \text{ large}), \\ t^r Q^{(r-1)}(t) &= O[(t^2/p^2)^r], \quad (t \rightarrow 0), \\ q^{(r)}(t) &= O(1/t^{r+1}), \quad (t \text{ large}), \\ t^r q^{(r)}(t) &= O(p^2/t^3), \quad (t \rightarrow 0). \end{aligned}$$

*Lemma 2 :*

If  $J(u) = \{1/\Gamma(h-\alpha+1)\} \int_u^\infty (t-u)^{h-\alpha} Q^{(h-1)}(t) dt$ , where  $\alpha \geq 2$  and  $h = [\alpha]$ , then

$$J(u) = O(u^{1-\alpha}), \quad (u \text{ large}).$$

*Lemma 3 :*

If  $J(u)$  is defined as in Lemma 2, then

$$\begin{aligned} (d/du) \{J(u)\} &= \{1/\Gamma(h-\alpha+1)\} \int_u^\infty (t-u)^{h-\alpha} Q^{(h)}(t) dt \\ &= O(u^{2h-\alpha+2}/p^{2h+2}), \quad (u \text{ small}). \end{aligned}$$

*Proof of Lemma 3 :*

Integrating by parts once, we have

$$\begin{aligned} J(u) &= \{1/\Gamma(h-\alpha+1)\} \int_u^\infty (t-u)^{h-\alpha} Q^{(h-1)}(t) dt \\ &= \{1/\Gamma(h-\alpha+2)\} \left[ \{Q^{(h-1)}(t) \cdot (t-u)^{h-\alpha+1}\}_u^\infty \right. \\ &\quad \left. - \int_u^\infty (t-u)^{h-\alpha+1} Q^{(h)}(t) dt \right] \\ &= \{1/\Gamma(h-\alpha+2)\} \left[ \lim_{t \rightarrow \infty} \{t^{h-\alpha+1}(1-u/t)^{h-\alpha+1} O(1/t^h)\} \right. \\ &\quad \left. - \int_u^\infty (t-u)^{h-\alpha+1} Q^{(h)}(t) dt \right] \\ &= -\{1/\Gamma(h-\alpha+2)\} \int_u^\infty (t-u)^{h-\alpha+1} Q^{(h)}(t) dt. \end{aligned}$$

Hence

$$(d/du)\{J(u)\} = \{1/\Gamma(h-\alpha+1)\} \int_u^\infty (t-u)^{h-\alpha} Q^{(h)}(t) dt.$$

*Lemma 4 :*

If  $J(u)$  is defined as in Lemma 2, we have

$$(d/du)^2\{J(u)\} = \{1/\Gamma(h-\alpha+1)\} \int_u^\infty (t-u)^{h-\alpha} Q^{(h+1)}(t) dt.$$

*Lemma 5 :*

$$\text{If } J^*(u) = \{(-1)^h/\Gamma\alpha\Gamma(h-\alpha+1)\} \int_u^\infty (t-u)^{(h-\alpha)} q^{(h)}(t) dt$$

where  $\alpha \geq 2$  and  $h = [\alpha]$ , then

$$\begin{aligned} J^*(u) &= O(u^{-\alpha}), \quad (u \text{ large}), \\ &= O(p^2 u^{-\alpha-2}), \quad (u \text{ small}). \end{aligned}$$

*Lemma 6 :*

If  $J^*(u)$  is defined as in Lemma 5, we have

$$(d/du)\{J^*(u)\} = \{(-1)^h/\Gamma\alpha\Gamma(h-\alpha+1)\} \int_u^\infty (t-u)^{h-\alpha} q^{(h+1)}(t) dt.$$

*Lemma 7 :*

$$\text{If } J'(u) = \{1/\Gamma(2-\alpha)\} \int_u^\infty (t-u)^{1-\alpha} Q^{(1)}(t) dt,$$

where  $1 \leq \alpha < 2$ , then

$$J'(u) = O(u^{-\alpha}), \quad (u \text{ large}).$$

*Lemma 8 :*

If  $J'(u)$  is defined as in Lemma 7, we have

$$(d/du)\{J'(u)\} = \{1/\Gamma(2-\alpha)\} \int_u^\infty (t-u)^{1-\alpha} Q^{(2)}(t) dt.$$

### 6.1. PROOF OF THEOREM 1.

Let

$$I = \frac{1}{2} \int_0^\infty \psi(t) Q(t) dt.$$

Then we have

$$\begin{aligned} I &= \frac{1}{4} \int_{-\infty}^\infty \psi(t) \cdot \{t/(p^2+t^2)\} dt \\ &= \frac{1}{4} \left[ \dots + \int_{-3\pi}^{-\pi} + \int_{-\pi}^{\pi} + \int_{\pi}^{3\pi} + \dots \right] \psi(t) \cdot \{t/(p^2+t^2)\} dt \\ &= \frac{1}{4} \left[ \int_{-\pi}^{\pi} \psi(t) \cdot \{t/(p^2+t^2)\} dt + \sum_{-\infty}^{\infty} \int_{(2k-1)\pi}^{(2k+1)\pi} \psi(t) \cdot \{t/(p^2+t^2)\} dt \right] \end{aligned}$$

$$\begin{aligned}
 &= \frac{1}{4} \left[ \int_{-\pi}^{\pi} \psi(t) \cdot \{t/(p^2+t^2)\} dt + \sum_{-\infty}^{\infty} \int_{-\pi}^{\pi} \psi(\tau+2k\pi) \frac{\tau+2k\pi}{p^2+(\tau+2k\pi)^2} d\tau \right] \\
 &= \frac{1}{4} \int_{-\pi}^{\pi} \left[ \sum_{-\infty}^{\infty} \frac{\tau+2k\pi}{p^2+(\tau+2k\pi)^2} + \frac{\tau}{p^2+\tau^2} \right] \cdot \psi(\tau) d\tau .
 \end{aligned}$$

Now from  $-\pi < \tau < \pi$ ,

$$\begin{aligned}
 &\sum_{-\infty}^{\infty} \frac{\tau+2k\pi}{p^2+(\tau+2k\pi)^2} + \frac{\tau}{p^2+\tau^2} \\
 &= \sum_{-\infty}^{\infty} \frac{1}{2i} \left[ \frac{1}{p-i(\tau+2k\pi)} - \frac{1}{p+i(\tau+2k\pi)} \right] + \frac{1}{2i} \left[ \frac{1}{p-i\tau} - \frac{1}{p+i\tau} \right] \\
 &= \frac{1}{2} \left[ \frac{1}{\tau+ip} + \sum_{-\infty}^{\infty} \frac{1}{(\tau+ip)+2k\pi} \right] + \frac{1}{2} \left[ \frac{1}{\tau-ip} + \sum_{-\infty}^{\infty} \frac{1}{(\tau-ip)+2k\pi} \right] \\
 &= \frac{1}{4} \left[ \frac{1}{\tau+ip} + 2 \left( \frac{\tau+ip}{2} \right) \sum_1^{\infty} \frac{1}{\left( \frac{\tau+ip}{2} \right)^2 - \pi^2 k^2} \right] \\
 &\qquad\qquad\qquad + \frac{1}{4} \left[ \frac{1}{\tau-ip} + 2 \left( \frac{\tau-ip}{2} \right) \sum_1^{\infty} \frac{1}{\left( \frac{\tau-ip}{2} \right)^2 - \pi^2 k^2} \right] \\
 &= \frac{1}{4} [\cot \frac{1}{2}(\tau+ip) + \cot \frac{1}{2}(\tau-ip)] \text{ [Titchmarsh, 1939.]} \\
 &= \frac{1}{2} [\sin \tau / (\cos ip - \cos \tau)] \\
 &= x \sin \tau / (1 - 2x \cos \tau + x^2) \\
 &= Q'(\tau).
 \end{aligned}$$

Thus finally

$$(6.11) \quad I = \frac{1}{4} \int_{-\pi}^{\pi} \psi(\tau) Q'(\tau) d\tau = \frac{1}{2} \int_0^{\pi} \psi(t) Q'(t) dt.$$

6.2. We have

$$\begin{aligned}
 V(x, \theta) &= \sum_1^{\infty} B_n(\theta) x^n \\
 &= (1/\pi) \int_{-\pi}^{\pi} f(\theta+t) Q'(t) dt \\
 &= (1/\pi) \int_0^{\pi} \psi(t) Q'(t) dt \\
 &= (1/\pi) \int_0^{\infty} \psi(t) Q(t) dt.
 \end{aligned}$$

Hence

$$\begin{aligned}
 V(x, \theta) - (1/\pi) \int_{\epsilon}^{\infty} \{\psi_{\alpha-1}(t)/t\} dt \\
 &= (1/\pi) \left[ \int_{\epsilon}^{\infty} \psi(t)Q(t) dt - \int_{\epsilon}^{\infty} \{\psi_{\alpha-1}(t)/t\} dt \right] \\
 (6.21) \quad &= (1/\pi)(I - I'), \text{ say.}
 \end{aligned}$$

Thus in order to establish Theorem 1, we have only to prove that if  $\psi_{\alpha}(t) = o(1)$ , then

$$I - I' = o(1).$$

6.3. We first consider the case in which  $\alpha \geq 2$ .

Integrating by parts  $(h-1)$  times, where  $h$  is the greatest integer not greater than  $\alpha$ , we have

$$\begin{aligned}
 I &= \int_0^{\infty} \psi(t)Q(t) dt \\
 &= \left[ \sum_{\rho=1}^{h-1} (-1)^{\rho-1} \Psi_{\rho}(t)Q^{(\rho-1)}(t) \right]_0^{\infty} + (-1)^{h-1} \int_0^{\infty} \Psi_{h-1}(t)Q^{(h-1)}(t) dt \\
 &= \left[ \sum_{\rho=1}^{h-1} o(t^{\rho}) \cdot O(t^{-\rho}) \right]_{t \rightarrow \infty} + (-1)^{h-1} I_1, \quad (\text{by Lemma 1}), \\
 &= o(1) + (-1)^{h-1} I_1.
 \end{aligned}$$

Now

$$\begin{aligned}
 I_1 &= \int_0^{\infty} \Psi_{h-1}(t) Q^{(h-1)}(t) dt \\
 &= \{1/\Gamma(h-\alpha+1)\} \int_0^{\infty} Q^{(h-1)}(t) dt \int_0^t (t-u)^{h-\alpha} d\Psi_{\alpha-1}(u) \\
 &= \{1/\Gamma(h-\alpha+1)\} \int_0^{\infty} d\Psi_{\alpha-1}(u) \int_u^{\infty} (t-u)^{h-\alpha} Q^{(h-1)}(t) dt \\
 &= \int_0^{\infty} J(u) d\Psi_{\alpha-1}(u).
 \end{aligned}$$

Integrating by parts once, we get

$$\begin{aligned}
 I_1 &= \left[ \Psi_{\alpha-1}(u) J(u) \right]_0^{\infty} - \int_0^{\infty} \Psi_{\alpha-1}(u) \cdot (d/du) \{J(u)\} du \\
 &= [o(u^{\alpha-1}) \cdot O(u^{1-\alpha})]_{u \rightarrow \infty} - I_2, \quad (\text{by Lemma 2}), \\
 &= o(1) - I_2.
 \end{aligned}$$

Now

$$\begin{aligned}
 I_2 &= \int_0^\infty \Psi_{\alpha-1}(u) \cdot (d/du) \{J(u)\} du \\
 &= (1/\Gamma\alpha) \int_0^\infty u^{\alpha-1} \psi_{\alpha-1}(u) \cdot (d/du) \{J(u)\} du \\
 &= (1/\Gamma\alpha) \left( \int_0^\epsilon + \int_\epsilon^\infty \right) u^{\alpha-1} \psi_{\alpha-1}(u) \cdot (d/du) \{J(u)\} du \\
 &= (1/\Gamma\alpha) (I_3 + I_4).
 \end{aligned}$$

We have now integrating once by parts

$$\begin{aligned}
 I_3 &= \int_0^\epsilon u^{\alpha-1} \psi_{\alpha-1}(u) \cdot (d/du) \{J(u)\} du \\
 &= [u^\alpha \psi_\alpha(u) \cdot (d/du) \{J(u)\}]_0^\epsilon - (\alpha-1) \int_0^\epsilon u^{\alpha-1} \psi_\alpha(u) (d/du) \{J(u)\} du \\
 &\quad - \int_0^\epsilon u^\alpha \psi_\alpha(u) \cdot (d/du)^2 \{J(u)\} du \\
 &= \epsilon^\alpha \psi_\alpha(\epsilon) \cdot O(\epsilon^{2h-\alpha+2}/p^{2h+2}) - (\alpha-1) I_5 - I_6, \text{ (by Lemma 3).} \\
 &= o(1) - (\alpha-1) I_5 - I_6, \text{ [since } \epsilon^2/p^2 = O(1), \text{ as } \epsilon \rightarrow 0. \text{]}
 \end{aligned}$$

Thus we have

$$(6.31) \quad I = o(1) + [(-1)^h/\pi\Gamma\alpha] I_4 - (\alpha-1) I_5 - I_6.$$

We first show that  $I_5 = o(1)$  and also  $I_6 = o(1)$ .

We have

$$\begin{aligned}
 I_5 &= \int_0^\epsilon u^{\alpha-1} \psi_\alpha(u) \cdot (d/du) \{J(u)\} du \\
 &= \{1/\Gamma(h-\alpha+1)\} \int_0^\epsilon u^{\alpha-1} \psi_\alpha(u) du \int_u^\infty (t-u)^{h-\alpha} Q^{(h)}(t) dt, \\
 &\hspace{25em} \text{(by Lemma 3),} \\
 &= \{1/\Gamma(h-\alpha+1)\} \int_\epsilon^\infty Q^{(h)}(t) dt \int_0^\epsilon u^{\alpha-1} \psi_\alpha(u) (t-u)^{h-\alpha} du \\
 &\quad + \{1/\Gamma(h-\alpha+1)\} \int_0^\epsilon u^{\alpha-1} \psi_\alpha(u) du \int_u^\epsilon (t-u)^{h-\alpha} Q^{(h)}(t) dt \\
 &= (I_{6,1} + I_{6,2}) / \Gamma(h-\alpha+1).
 \end{aligned}$$

Now

$$\begin{aligned}
 I_{5,1} &= \int_{\epsilon}^{\infty} Q^{(h)}(t) dt \int_0^{\epsilon} u^{\alpha-1} \psi_{\alpha}(u) (t-u)^{h-\alpha} du \\
 &= O \left[ \int_{\epsilon}^{\infty} t^{-h-1} dt \int_0^{\epsilon} o(1) \cdot u^{\alpha-1} (t-u)^{h-\alpha} du \right], \text{ (by Lemma 1),} \\
 &= o(1) \cdot O \left[ \int_{\epsilon}^{\infty} t^{-h-1} \epsilon^{\alpha-1} dt \int_0^{\epsilon} (t-u)^{h-\alpha} du \right] \\
 &= o(1) \cdot O \left[ \int_{\epsilon}^{\infty} \epsilon^{\alpha-1} t^{-h-1} t^{h-\alpha+1} dt \right] \\
 &= o(1) \cdot O \left[ \epsilon^{\alpha-1} / t^{\alpha-1} \right]_{\epsilon}^{\infty} \\
 &= o(1).
 \end{aligned}$$

$$\begin{aligned}
 I_{5,2} &= \int_0^{\epsilon} Q^{(h)}(t) dt \int_0^t u^{\alpha-1} \psi_{\alpha}(u) (t-u)^{h-\alpha} du \\
 &= o(1) \cdot \int_0^{\epsilon} t^h Q^{(h)}(t) dt \int_0^1 x^{\alpha-1} (1-x)^{h-\alpha} dx \\
 &= o(1) \cdot O \left[ \int_0^{\epsilon} (t^{2h+1} / p^{2h+2}) dt \right] \\
 &= o(1) \cdot O(\epsilon^{2h+2} / p^{2h+2}), \left[ \text{since } \int_0^1 x^{\alpha-1} (1-x)^{h-\alpha} dx < \infty \right], \\
 &= o(1).
 \end{aligned}$$

Next

$$\begin{aligned}
 I_6 &= \int_0^{\epsilon} u^{\alpha} \psi_{\alpha}(u) (d/du)^2 \{J(u)\} du \\
 &= \{1/\Gamma(h-\alpha+1)\} \int_0^{\epsilon} u^{\alpha} \psi_{\alpha}(u) du \int_u^{\infty} (t-u)^{h-\alpha} Q^{(h+1)}(t) dt, \\
 &\hspace{25em} \text{(by Lemma 4),} \\
 &= \{1/\Gamma(h-\alpha+1)\} \left[ \int_{\epsilon}^{\infty} Q^{(h+1)}(t) dt \int_0^{\epsilon} u^{\alpha} \psi_{\alpha}(u) (t-u)^{h-\alpha} du \right. \\
 &\hspace{15em} \left. + \int_0^{\epsilon} u^{\alpha} \psi_{\alpha}(u) du \int_u^{\epsilon} (t-u)^{h-\alpha} Q^{(h+1)}(t) dt \right] \\
 &= (I_{6,1} + I_{6,2}) / \Gamma(h-\alpha+1).
 \end{aligned}$$



Now

$$\begin{aligned}
 I_{6,1} &= \int_{\epsilon}^{\infty} Q^{(h+1)}(t) dt \int_0^{\epsilon} u^{\alpha} \psi_{\alpha}(u)(t-u)^{h-\alpha} du \\
 &= O \left[ \int_{\epsilon}^{\infty} t^{-h-2} dt \int_0^{\epsilon} o(1) \cdot u^{\alpha} (t-u)^{h-\alpha} du \right], \text{ (by Lemma 1),} \\
 &= o(1) \cdot O \left[ \int_{\epsilon}^{\infty} t^{-h-2} \epsilon^{\alpha} dt \int_0^{\epsilon} (t-u)^{h-\alpha} du \right] \\
 &= o(1) \cdot O \left[ \int_{\epsilon}^{\infty} \epsilon^{\alpha} t^{-h-2} t^{h-\alpha+1} dt \right] \\
 &= o(1) \cdot O \left[ \int_{\epsilon}^{\infty} (\epsilon^{\alpha}/t^{\alpha+1}) dt \right] \\
 &= o(1) \cdot O[\epsilon^{\alpha}/t^{\alpha}]_{\epsilon}^{\infty} \\
 &= o(1) . \\
 I_{6,2} &= \int_0^{\epsilon} Q^{(h+1)}(t) dt \int_0^t u^{\alpha} \psi_{\alpha}(u)(t-u)^{h-\alpha} du \\
 &= o(1) \cdot \int_0^{\epsilon} t^{h+1} Q^{(h+1)}(t) dt \int_0^1 x^{\alpha}(1-x)^{h-\alpha} dx \\
 &= o(1) \cdot O \left[ \int_0^{\epsilon} (t^{2h+3}/p^{2h+4}) dt \right], \text{ (by Lemma 1),} \\
 &= o(1) \cdot O(\epsilon^{2h+4}/p^{2h+4}) \\
 &= o(1) .
 \end{aligned}$$

Hence now

$$(6.32) \quad I = o(1) + \{(-1)^h/\pi\Gamma\alpha\}I_4,$$

so that combining (6.21) and (6.32), we have

$$\begin{aligned}
 V(x,\theta) - (1/\pi) \int_{\epsilon}^{\infty} \{\psi_{\alpha-1}(u)/u\} du \\
 &= o(1) + \{(-1)^h/\pi\Gamma\alpha\} \cdot I_4 - (1/\pi)I' \\
 &= o(1) + K \cdot \int_{\epsilon}^{\infty} u^{\alpha-1} \psi_{\alpha-1}(u) du \int_u^{\infty} (t-u)^{h-\alpha} Q^{(h)}(t) dt - (1/\pi)I', \\
 &\hspace{20em} [\text{where } K = \{(-1)^h/\pi\Gamma\alpha\Gamma(h-\alpha+1)\}], \\
 &= o(1) + K \cdot \left[ \int_{\epsilon}^{\infty} u^{\alpha-1} \psi_{\alpha-1}(u) du \int_u^{\infty} (t-u)^{h-\alpha} \left\{ \frac{(-1)^h \Gamma(h+1)}{t^{h+1}} - q^{(h)}(t) \right\} dt \right] - (1/\pi)I'
 \end{aligned}$$

$$\begin{aligned}
 &= o(1) + K \cdot (-1)^h \Gamma(h+1) \cdot \int_{\epsilon}^{\infty} u^{\alpha-1} \psi_{\alpha-1}(u) du \int_u^{\infty} (t-u)^{h-\alpha} t^{-h-1} dt \\
 &\hspace{25em} -(1/\pi)I' - I_7 \\
 &= o(1) + 2(-1)^h K \cdot \Gamma(h+1) \int_{\epsilon}^{\infty} u^{\alpha-1} \psi_{\alpha-1}(u) du \cdot \int_0^{\pi/2} u^{-\alpha} (\sin \theta)^{2\alpha-1} \\
 &\hspace{25em} \times (\cos \theta)^{2h-2\alpha+1} d\theta - I_7 - (1/\pi)I' \\
 &= o(1) + (1/\pi) \int_{\epsilon}^{\infty} \{\psi_{\alpha-1}(u)/u\} du - I_7 - (1/\pi)I' \\
 &= o(1) - I_7.
 \end{aligned}$$

It now remains to show that  $I_7 = o(1)$ .

Now

$$\begin{aligned}
 I_7 &= K \cdot \left[ \int_{\epsilon}^{\infty} u^{\alpha-1} \psi_{\alpha-1}(u) du \int_u^{\infty} (t-u)^{h-\alpha} q^{(h)}(t) dt \right] \\
 &= \int_{\epsilon}^{\infty} u^{\alpha-1} \psi_{\alpha-1}(u) J^*(u) du.
 \end{aligned}$$

Integrating by parts once, we get

$$\begin{aligned}
 I_7 &= [u^{\alpha-1} J^*(u) \cdot u \psi_{\alpha}(u)]_{\epsilon}^{\infty} \\
 &\quad - \int_{\epsilon}^{\infty} u \psi_{\alpha}(u) \left[ (\alpha-1) u^{\alpha-2} J^*(u) + u^{\alpha-1} (d/du) \{J^*(u)\} du \right] \\
 &= O \left[ u^{\alpha} \psi_{\alpha}(u) \cdot u^{-\alpha} \right] + O \left[ u^{\alpha} \psi_{\alpha}(u) \cdot (p^2 u^{-\alpha-2}) \right]_{u=\epsilon}^{\infty} \\
 &\hspace{25em} - (\alpha-1) I_8 - I_9 \\
 &= o(1) - (\alpha-1) I_8 - I_9.
 \end{aligned}$$

We now proceed to show that  $I_8 = o(1)$  and  $I_9 = o(1)$ .

Now

$$\begin{aligned}
 I_8 &= \int_{\epsilon}^{\infty} u^{\alpha-1} \psi_{\alpha}(u) J^*(u) du \\
 &= K \cdot \int_{\epsilon}^{\infty} u^{\alpha-1} \psi_{\alpha}(u) du \int_u^{\infty} (t-u)^{h-\alpha} q^{(h)}(t) dt \\
 &= O \left[ p^2 \int_{\epsilon}^{\infty} u^{\alpha-1} \psi_{\alpha}(u) \cdot u^{-\alpha-2} du \int_0^1 x^{\alpha+1} (1-x)^{h-\alpha} dx \right], \\
 &\hspace{25em} \text{(by Lemma 1),}
 \end{aligned}$$

$$\begin{aligned}
 &= O \left[ p^2 \left( \int_{\epsilon}^{\delta} + \int_{\delta}^{\infty} \right) u^{-3} |\psi_{\alpha}(u)| du \right] \\
 &= O(I_{8,1} + I_{8,2}).
 \end{aligned}$$

$$\begin{aligned}
 I_{8,1} &= p^2 \int_{\epsilon}^{\delta} u^{-3} |\psi_{\alpha}(u)| du \\
 &= o \left( p^2 \int_{\epsilon}^{\delta} u^{-3} du \right) \\
 &= o(p^2/\epsilon^2) \\
 &= o(1), \quad \text{since } p = o(1).
 \end{aligned}$$

$$\begin{aligned}
 I_{8,2} &= p^2 \int_{\delta}^{\infty} u^{-3} |\psi_{\alpha}(u)| du \\
 &= p^2 \int_{\delta}^{\infty} u^{-3} \cdot O(1) \cdot du \\
 &= O \left[ p^2 \int_{\delta}^{\infty} u^{-3} du \right] \\
 &= o(1).
 \end{aligned}$$

Next

$$\begin{aligned}
 I_9 &= \int_{\epsilon}^{\infty} u^{\alpha} \psi_{\alpha}(u) (d/du) \{ J^*(u) \} du \\
 &= K \cdot \int_{\epsilon}^{\infty} u^{\alpha} \psi_{\alpha}(u) du \int_u^{\infty} (t-u)^{h-\alpha} q^{(h+1)}(t) dt \\
 &= O \left[ p^2 \int_{\epsilon}^{\infty} u^{\alpha} |\psi_{\alpha}(u)| u^{-\alpha-3} du \int_0^1 x^{\alpha+2} (1-x)^{h-\alpha} dx \right], \\
 & \hspace{25em} \text{(by Lemma 1),} \\
 &= O \left[ p^2 \left( \int_{\epsilon}^{\delta} + \int_{\delta}^{\infty} \right) u^{-3} |\psi_{\alpha}(u)| du \right] \\
 &= O(I_{9,1} + I_{9,2}).
 \end{aligned}$$

Now

$$\begin{aligned}
 I_{9,1} &= p^2 \int_{\epsilon}^{\delta} u^{-3} |\psi_{\alpha}(u)| du \\
 &= o \left( p^2 \int_{\epsilon}^{\delta} u^{-3} du \right) \\
 &= o(p^2/\epsilon^2) \\
 &= o(1).
 \end{aligned}$$

$$\begin{aligned}
I_{0,2} &= p^2 \int_{\delta}^{\infty} u^{-3} |\psi_{\alpha}(u)| du \\
&= p^2 \int_{\delta}^{\infty} u^{-3} \cdot O(1) \cdot du \\
&= O\left(p^2 \int_{\delta}^{\infty} u^{-3} du\right) \\
&= o(1).
\end{aligned}$$

This completes the proof of Theorem 1 when  $\alpha \geq 2$ .

6.4. We now proceed to give the proof for the case in which  $1 \leq \alpha < 2$ .

From (6.21) we know that in order to establish the theorem we have to prove that if  $\psi_{\alpha}(t) = o(1)$ ,  $1 \leq \alpha < 2$ , then

$$I - I' = o(1).$$

Integrating once by parts, we get

$$\begin{aligned}
I &= \int_0^{\infty} \psi(t) Q(t) dt \\
&= [Q(t)\Psi_1(t)]_0^{\infty} - \int_0^{\infty} \Psi_1(t)Q^{(1)}(t) dt \\
&= o(1) - \{1/\Gamma(2-\alpha)\} \int_0^{\infty} Q^{(1)}(t) dt \int_0^t (t-u)^{1-\alpha} d\Psi_{\alpha}(u) \\
&= o(1) - \int_0^{\infty} J'(u)d\Psi_{\alpha}(u) \\
&= o(1) + I_1.
\end{aligned}$$

Now integrating by parts once we get

$$\begin{aligned}
I_1 &= - \int_0^{\infty} J'(u)d\Psi_{\alpha}(u) \\
&= - [J'(u)\Psi_{\alpha}(u)]_0^{\infty} + \int_0^{\infty} \Psi_{\alpha}(u)(d/du) \{J'(u)\} du \\
&= [O(u^{-\alpha}) \cdot o(u^{\alpha})]_{u \rightarrow \infty} + I_2 \\
&= o(1) + I_2.
\end{aligned}$$

Now

$$\begin{aligned}
I_2 &= \int_0^{\infty} \Psi_{\alpha}(u)(d/du) \{J'(u)\} du \\
&= \{1/\Gamma(\alpha+1)\} \left( \int_0^{\epsilon} + \int_{\epsilon}^{\infty} \right) u^{\alpha} \psi_{\alpha}(u) \cdot (d/du) \{J'(u)\} du \\
&= (I_3 + I_4)/\Gamma(\alpha+1).
\end{aligned}$$

We now show that  $I_3 = o(1)$ .

Now

$$\begin{aligned}
 I_3 &= \int_0^\epsilon u^\alpha \psi_\alpha(u) (d/du) \{J'(u)\} du \\
 &= \{1/\Gamma(2-\alpha)\} \int_0^\epsilon u^\alpha \psi_\alpha(u) du \int_u^\infty (t-u)^{1-\alpha} Q^{(2)}(t) dt \\
 &= \{1/\Gamma(2-\alpha)\} \left[ \int_0^\epsilon u^\alpha \psi_\alpha(u) du \int_u^\epsilon (t-u)^{1-\alpha} Q^{(2)}(t) dt \right. \\
 &\quad \left. + \int_\epsilon^\infty Q^{(2)}(t) dt \int_0^\epsilon u^\alpha \psi_\alpha(u) (t-u)^{1-\alpha} du \right] \\
 &= (I_{3,1} + I_{3,2})/\Gamma(2-\alpha). \\
 I_{3,1} &= \int_0^\epsilon u^\alpha \psi_\alpha(u) du \int_u^\epsilon (t-u)^{1-\alpha} Q^{(2)}(t) dt \\
 &= \int_0^\epsilon Q^{(2)}(t) dt \int_0^t u^\alpha \psi_\alpha(u) (t-u)^{1-\alpha} du \\
 &= o(1) \cdot \int_0^\epsilon t^2 Q^{(2)}(t) dt \int_0^1 x^\alpha (1-x)^{1-\alpha} dx \\
 &= o(1) \cdot O \left[ \int_0^\epsilon (t^5/p^6) dt \right] \\
 &= o(1) \cdot O(\epsilon^6/p^6) \\
 &= o(1). \\
 I_{3,2} &= \int_\epsilon^\infty Q^{(2)}(t) dt \int_0^\gamma u^\alpha \psi_\alpha(u) (t-u)^{1-\alpha} du \\
 &= o(1) \cdot O \left[ \int_\epsilon^\infty t^{-3\epsilon^\alpha} dt \int_0^\epsilon (t-u)^{1-\alpha} du \right] \\
 &= o(1) \cdot O \left[ \int_\epsilon^\infty \epsilon^\alpha t^{-3t^{2-\alpha}} dt \right] \\
 &= o(1) \cdot O \left[ \int_\epsilon^\infty (\epsilon^\alpha/t^{\alpha+1}) dt \right] \\
 &= o(1) \cdot O[(\epsilon^\alpha/t^\alpha)]_\epsilon^\infty \\
 &= o(1).
 \end{aligned}$$

(6.41) Thus

$$I = o(1) + [1/\pi\Gamma(\alpha + 1)] \cdot I_4$$

Now we have

$$\begin{aligned} V(x, \theta) &= (1/\pi) \int_{\epsilon}^{\infty} \{\psi_{\alpha-1}(u)/u\} du \\ &= o(1) + [1/\pi\Gamma(\alpha + 1)] \cdot I_4 - (1/\pi) I' \\ &= o(1) + K \cdot \int_{\epsilon}^{\infty} u^{\alpha} \psi_{\alpha}(u) du \int_u^{\infty} (t-u)^{1-\alpha} Q^{(2)}(t) dt - (1/\pi) I', \end{aligned}$$

where

$$\begin{aligned} & [K = \{1/\pi\Gamma(\alpha + 1)\Gamma(2-\alpha)\}], \\ &= o(1) + K \cdot \int_{\epsilon}^{\infty} u^{\alpha} \psi_{\alpha}(u) du \int_u^{\infty} (t-u)^{1-\alpha} \{ (2/t^3) - q^{(2)}(t) \} dt - (1/\pi) I' \\ &= o(1) + 2K \int_{\epsilon}^{\infty} u^{\alpha} \psi_{\alpha}(u) du \int_u^{\infty} (t-u)^{1-\alpha} t^{-3} dt - I_5 - (1/\pi) I' \\ &= o(1) + 4K \int_{\epsilon}^{\infty} u^{-1} \psi_{\alpha}(u) du \int_{\pi/2}^{\infty} (\sin \theta)^{3-2\alpha} (\cos \theta)^{2\alpha+1} d\theta \\ & \qquad \qquad \qquad - I_5 - (1/\pi) I' \\ &= o(1) + (1/\pi) \int_{\epsilon}^{\infty} u^{-1} \psi_{\alpha}(u) du - I_5 - (1/\pi) \int_{\epsilon}^{\infty} u^{-1} \psi_{\alpha-1}(u) du \\ &= o(1) + (1/\pi) \int_{\epsilon}^{\infty} u^{-1} \psi_{\alpha}(u) du - I_5 - (1/\pi) \left[ \psi_{\alpha}(u) \right]_{\epsilon}^{\infty} \\ & \qquad \qquad \qquad - (1/\pi) \int_{\epsilon}^{\infty} u^{-1} \psi_{\alpha}(u) du \\ &= o(1) - I_5. \end{aligned}$$

It now remains to show that  $I_5 = o(1)$ .

$$\begin{aligned} I_5 &= K \int_{\epsilon}^{\infty} u^{\alpha} \psi_{\alpha}(u) du \int_u^{\infty} (t-u)^{1-\alpha} Q^{(2)}(t) dt \\ &= O \left[ p^2 \int_{\epsilon}^{\infty} u^{\alpha} |\psi_{\alpha}(u)| \cdot u^{-\alpha-3} du \int_0^1 x^{\alpha+2} (1-x)^{1-\alpha} dx \right] \\ &= O \left[ p^2 \left( \int_{\epsilon}^{\delta} + \int_{\delta}^{\infty} \right) u^{-3} |\psi_{\alpha}(u)| du \right] \\ &= O(I_{5, 1} + I_{5, 2}). \end{aligned}$$

Now

$$\begin{aligned} I_{5,1} &= p^2 \int_{\epsilon}^{\delta} u^{-3} |\psi_{\alpha}(u)| du \\ &= o\left(p^2 \int_{\epsilon}^{\delta} u^{-3} du\right) \\ &= o(1). \end{aligned}$$

$$\begin{aligned} I_{5,2} &= p^2 \int_{\delta}^{\infty} u^{-3} |\psi_{\alpha}(u)| du \\ &= O\left(p^2 \int_{\delta}^{\infty} u^{-3} du\right) \\ &= o(1). \end{aligned}$$

This completes the proof of Theorem 1.

7. *Proof of Theorem 2*—Theorem 2 can be proved by arguments parallel to those used in the proof of Theorem 1 with the only difference that we have to use the condition

$$\psi_{\alpha}(t) = O(1)$$

instead of

$$\psi_{\alpha}(t) = o(1)$$

so long as  $\alpha \geq 1$ . But when  $0 < \alpha < 1$ , we prove the theorem as follows :

$$\begin{aligned} I &= \int_0^{\infty} \psi(t) Q(t) dt \\ &= \{1/\Gamma(1-\alpha)\} \int_0^{\infty} Q(t) dt \int_0^t (t-u)^{-\alpha} d\Psi_{\alpha}(u) \\ &= \{1/\Gamma(1-\alpha)\} \int_0^{\infty} d\Psi_{\alpha}(u) \int_u^{\infty} (t-u)^{-\alpha} Q(t) dt \\ &= \int_0^{\infty} J_1(u) d\Psi_{\alpha}(u), \end{aligned}$$

where

$$J_1(u) = \frac{1}{\Gamma(1-\alpha)} \int_u^{\infty} (t-u)^{-\alpha} Q(t) dt = O(u^{-\alpha}), \text{ as } u \rightarrow \infty.$$

Integrating by parts once, we get

$$I = [\Psi_{\alpha}(u)J_1(u)]_0^{\infty} - \int_0^{\infty} \Psi_{\alpha}(u) (d/du)\{J_1(u)\} du,$$

where

$$(d/du)\{J_1(u)\} = \frac{1}{\Gamma(1-\alpha)} \int_u^\infty (t-u)^{-\alpha} Q^{(1)}(t) dt.$$

$$\begin{aligned} &= [O(u^\alpha) \cdot O(u^{-\alpha})]_{u \rightarrow \infty} - I_1 \\ &= O(1) - I_1. \end{aligned}$$

Now

$$\begin{aligned} I_1 &= \int_0^\infty \Psi_\alpha(u) (d/du)\{J_1(u)\} du \\ &= \{1/\Gamma(\alpha+1)\} \left( \int_0^\epsilon + \int_\epsilon^\infty \right) u^\alpha \psi_\alpha(u) (d/du)\{J_1(u)\} du \\ &= (I_2 + I_3)/(\alpha+1). \end{aligned}$$

Now

$$\begin{aligned} I_2 &= \int_0^\epsilon u^\alpha \psi_\alpha(u) (d/du)\{J_1(u)\} du \\ &= \{1/\Gamma(1-\alpha)\} \left[ \int_0^\epsilon u^\alpha \psi_\alpha(u) du \int_u^\infty (t-u)^{-\alpha} Q^{(1)}(t) dt \right] \\ &= \frac{1}{\Gamma(1-\alpha)} \left[ \int_0^\epsilon Q^{(1)}(t) dt \int_0^t u^\alpha \cdot \psi_\alpha(u) \cdot (t-u)^{-\alpha} du \right. \\ &\quad \left. + \int_\epsilon^\infty Q^{(1)}(t) dt \int_0^\epsilon u^\alpha \psi_\alpha(u) \cdot (t-u)^{-\alpha} du \right] \\ &= (I_{2,1} + I_{2,2})/\Gamma(1-\alpha) \\ I_{2,1} &= \int_0^\epsilon Q^{(1)}(t) dt \int_0^t u^\alpha \psi_\alpha(u) (t-u)^{-\alpha} du \\ &= O(1) \cdot \int_0^\epsilon t Q^{(1)}(t) dt \int_0^1 x^\alpha (1-x)^{-\alpha} dx \\ &= O(1) \cdot O \left[ \int_0^\epsilon (t^3/p^4) dt \right] \\ &= O(1). \\ I_{2,2} &= \int_\epsilon^\infty Q^{(1)}(t) dt \int_0^\epsilon u^\alpha \psi_\alpha(u) (t-u)^{-\alpha} du \\ &= O(1) \cdot O \left[ \int_\epsilon^\infty t^{-2} \epsilon^\alpha dt \int_0^\epsilon (t-u)^{-\alpha} du \right] \\ &= O(1) \cdot O \left[ \int_\epsilon^\infty \epsilon^\alpha t^{-2} t^{1-\alpha} dt \right] \\ &= O(1) \cdot O \left[ (\epsilon^\alpha/t^\alpha) \right]_\epsilon^\infty \\ &= O(1). \end{aligned}$$



(7.1) Thus  $I = O(1) - \{1/\pi\Gamma(\alpha+1)\}I_3$

so that

$$\begin{aligned} V(x, \theta) &= (1/\pi) \int_{\epsilon}^{\infty} \{\psi_{\alpha-1}(u)/u\} du \\ &= O(1) - [1/\pi\Gamma(\alpha+1)]I_3 - (1/\pi) I' \\ &= O(1) - \left[1/\pi\Gamma(\alpha+1)\right] \int_{\epsilon}^{\infty} u^{\alpha}\psi_{\alpha}(u) du \int_u^{\infty} (t-u)^{-\alpha}Q^{(1)}(t) dt - (1/\pi)I' \\ &= O(1) - K \int_{\epsilon}^{\infty} u^{\alpha}\psi_{\alpha}(u) du \int_u^{\infty} (t-u)^{-\alpha}\{(-1/t^2) - q^{(1)}(t)\} dt - (1/\pi)I', \\ &\qquad\qquad\qquad \text{where } K = \{1/\pi\Gamma(1-\alpha)\Gamma(1+\alpha)\} \\ &= O(1) + K \int_{\epsilon}^{\infty} u^{\alpha}\psi_{\alpha}(u) du \int_u^{\infty} (t-u)^{-\alpha} \cdot t^{-2} dt + I_4 - (1/\pi)I' \\ &= O(1) + 2K \int_{\epsilon}^{\infty} u^{-1}\psi_{\alpha}(u) du \int_0^{\pi/2} (\sin \theta)^{1-2\alpha}(\cos \theta)^{1+2\alpha}d\theta + I_4 - (1/\pi)I' \\ &= O(1) + (1/\pi) \int_{\epsilon}^{\infty} u^{-1}\psi_{\alpha}(u) du + I_4 - (1/\pi) \int_{\epsilon}^{\infty} u^{-1}\psi_{\alpha-1}(u) du \\ &= O(1) + I_4. \end{aligned}$$

It now remains to show that  $I_4 = O(1)$ .

Now

$$\begin{aligned} I_4 &= K \int_{\epsilon}^{\infty} u^{\alpha}\psi_{\alpha}(u) du \int_u^{\infty} (t-u)^{-\alpha}q^{(1)}(t) dt \\ &= O \left[ p^2 \int_{\epsilon}^{\infty} u^{\alpha} |\psi_{\alpha}(u)| \cdot u^{-\alpha-3} du \int_0^1 x^{\alpha+2}(1-x)^{-\alpha} dx \right] \\ &= O \left[ p^2 \left( \int_{\epsilon}^{\delta} + \int_{\delta}^{\infty} \right) u^{-3} |\psi_{\alpha}(u)| du \right] \\ &= O [I_{4, 1} + I_{4, 2}]. \\ I_{4, 1} &= p^2 \int_{\epsilon}^{\delta} u^{-3} |\psi_{\alpha}(u)| du \\ &= O(1). \\ I_{4, 2} &= p^2 \int_{\delta}^{\infty} u^{-3} |\psi_{\alpha}(u)| du \\ &= O(1). \end{aligned}$$

This completes the proof of Theorem 2.

It is obvious that when  $\alpha$  ( $\geq 1$ ) is a positive integer, i.e. when  $h = \alpha$ , there will arise an automatic simplification in the proof given above; a straightforward treatment of this case has already been given in the author's earlier paper [1953].

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