

DEFINITE INTEGRALS INVOLVING SELF-RECIPROCAL FUNCTIONS*

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In this paper, some theorems have been given about the functions such that if their $\chi_{\mu, k, m}$ -transform is $F(x)/\sqrt{x}$, their $\chi_{\nu, k, m}$ -transform is $F\left(\frac{1}{x}\right)/\sqrt{x}$. The nature of such functions has been discussed and it has been shown that definite integrals, which involve self-reciprocal functions, satisfy a particular type of relation. With the help of some of these results, self-reciprocal functions have been derived as solutions of integral equations.

§ 1. A generalization of the Hankel transform

$$g(x) = \int_0^{\infty} \sqrt{xy} J_{\nu}(xy) f(y) dy \quad \dots \quad \dots \quad \dots \quad (1.1)$$

has been introduced by Narain (1956-57, p. 270) in the form

$$g(x) = 2^{-\nu} \int_0^{\infty} (xy)^{\nu+\frac{1}{2}} \chi_{\nu, k, m}(x^2y^2/4) f(y) dy \quad \dots \quad \dots \quad (1.2)$$

where

$$\chi_{\nu, k, m}(x) = x^{-\nu} G_{2, 4}^{2, 1} \left(x \left| \begin{matrix} k-m-\frac{1}{2}, \nu-k+m+\frac{1}{2} \\ \nu, \nu+2m, -2m, 0 \end{matrix} \right. \right).$$

He calls $g(x)$ in (1.2) to be $\chi_{\nu, k, m}$ -transform of $f(x)$ and has proved in the same paper that this is a transform reciprocal like the Hankel transform. In case $g(x) \equiv f(x)$ we may call it to be self-reciprocal in the $\chi_{\nu, k, m}$ -transform and we shall say that $f(x)$ is $R_{\nu}(k, m)$.

When $k+m = \frac{1}{2}$, (1.2) reduces to (1.1).

Yet another generalization of (1.1) has been given by Narain (1959, p. 298) as:

$$g(x) = \sqrt{2} \int_0^{\infty} G_{2p, 2q}^{q, p} \left(\frac{x^2y^2}{4} \left| \begin{matrix} a_1, \dots, a_p, \frac{1}{2}-a_1, \dots, \frac{1}{2}-a_p \\ b_1, \dots, b_q, \frac{1}{2}-b_1, \dots, \frac{1}{2}-b_q \end{matrix} \right. \right) f(y) dy \quad \dots \quad (1.3)$$

On putting $p = 0$, $q = 1$ and $b_1 = \frac{1}{4} + \frac{1}{2}\nu$ (1.3) reduces to (1.1).

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In this paper, we have given some theorems about functions such that if their $\chi_{\mu, k, m}$ -transform is $x^{-\frac{1}{2}}F(x)$, their $\chi_{\nu, k, m}$ -transform is $x^{-\frac{1}{2}}F\left(\frac{1}{x}\right)$. We have found out the nature of such functions and have also shown that definite integrals, which involve self-reciprocal functions, satisfy a particular type of relation. With the help of some of the results, self-reciprocal functions have been derived as solutions of the integral equation.

§ 2. THEOREM 1. If $f(x)$ is $R_{\nu}(k, m)$, i.e. a solution of the integral equation

$$f(x) = \sqrt{2} \int_0^{\infty} G_{2,4}^{2,1} \left(\frac{x^2 t^2}{4} \left| \begin{matrix} k-m-\frac{1}{4}-\frac{\nu}{4}, \frac{\nu}{2}-k+m+\frac{3}{4} \\ \frac{\nu}{2}+\frac{1}{4}, \frac{\nu}{2}+\frac{1}{4}+2m, \frac{1}{4}-\frac{\nu}{2}, \frac{1}{4}-\frac{\nu}{2}-2m \end{matrix} \right. \right) f(t) dt$$

and if

$$g(x) = \int_a^{1/a} t^{-\frac{1}{2}} F(t) f(xt) dt$$

where a is positive or zero and $F(t)$ is any function such that

$$F(t) = F\left(\frac{1}{t}\right),$$

then $g(x)$ is also $R_{\nu}(k, m)$, i.e. is a solution of the integral equation, provided the integrals involved are absolutely and uniformly convergent.

PROOF: We have

$$\begin{aligned} & \sqrt{2} \int_0^{\infty} G_{2,4}^{2,1} \left(\frac{x^2 t^2}{4} \left| \begin{matrix} k-m-\frac{1}{4}-\frac{\nu}{2}, \frac{\nu}{2}-k+m+\frac{3}{4} \\ \frac{\nu}{2}+\frac{1}{4}, \frac{\nu}{2}+\frac{1}{4}+2m, \frac{1}{4}-\frac{\nu}{2}, \frac{1}{4}-\frac{\nu}{2}-2m \end{matrix} \right. \right) g(t) dt \\ &= \sqrt{2} \int_0^{\infty} G_{2,4}^{2,1} \left(\frac{x^2 t^2}{4} \left| \begin{matrix} k-m-\frac{1}{4}-\frac{\nu}{2}, \frac{\nu}{2}-k+m+\frac{3}{4} \\ \frac{\nu}{2}+\frac{1}{4}, \frac{\nu}{2}+\frac{1}{4}+2m, \frac{1}{4}-\frac{\nu}{2}, \frac{1}{4}-\frac{\nu}{2}-2m \end{matrix} \right. \right) dt \int_a^{1/a} u^{-\frac{1}{2}} F(u) f(tu) du \\ &= \sqrt{2} \int_a^{1/a} u^{-\frac{1}{2}} F(u) du \int_0^{\infty} G_{2,4}^{2,1} \left(\frac{x^2 t^2}{4} \left| \begin{matrix} k-m-\frac{1}{4}-\frac{\nu}{2}, \frac{\nu}{2}-k+m+\frac{3}{4} \\ \frac{\nu}{2}+\frac{1}{4}, \frac{\nu}{2}+\frac{1}{4}+2m, \frac{1}{4}-\frac{\nu}{2}, \frac{1}{4}-\frac{\nu}{2}-2m \end{matrix} \right. \right) f(tu) dt \\ &= \sqrt{2} \int_a^{1/a} u^{-\frac{1}{2}} F(u) du \int_0^{\infty} G_{2,4}^{2,1} \left(\frac{x^2 t^2}{4u^2} \left| \begin{matrix} k-m-\frac{1}{4}-\frac{\nu}{2}, \frac{\nu}{2}-k+m+\frac{3}{4} \\ \frac{\nu}{2}+\frac{1}{4}, \frac{\nu}{2}+\frac{1}{4}+2m, \frac{1}{4}-\frac{\nu}{2}, \frac{1}{4}-\frac{\nu}{2}-2m \end{matrix} \right. \right) f(t) dt \\ &= \int_a^{1/a} u^{-\frac{1}{2}} F(u) f\left(\frac{x}{u}\right) du = \int_a^{1/a} t^{-\frac{1}{2}} F\left(\frac{1}{t}\right) f(xt) dt \\ &= g(x), \end{aligned}$$

since $F(t) = F\left(\frac{1}{t}\right)$.

When $a = 0$, the theorem can be expressed in the form that

$$g(x) = \int_0^\infty (xu)^{-\frac{1}{2}} F(u/x) f(u) du$$

is self-reciprocal in $X_{\nu, k, m}$ -transform.

The most obvious example is obtained by taking

$$F(t) = (t^\alpha + t^{-\alpha})^{-\lambda},$$

in which case we get the self-reciprocal function in $X_{\nu, k, m}$ -transform

$$\int_0^\infty \frac{(xu)^{\alpha\lambda - \frac{1}{2}}}{(u^{2\alpha} + x^{2\alpha})^\lambda} f(u) du.$$

The argument can also be generalized when $a = 0$. The corresponding result is then:

THEOREM 2. If $f(x)$ is $R_\nu(k, m)$, i.e. a solution of the integral equation

$$f(x) = \sqrt{2} \int_0^\infty G_{2,4}^{2,1} \left(\frac{x^2 t^2}{4} \left| \begin{matrix} k-m-\frac{1}{4}-\frac{\nu}{2}, \frac{\nu}{2}-k+m+\frac{3}{4} \\ \frac{\nu}{2}+\frac{1}{4}, \frac{\nu}{2}+\frac{1}{4}+2m, \frac{1}{4}-\frac{\nu}{2}, \frac{1}{4}-\frac{\nu}{2}-2m \end{matrix} \right. \right) f(t) dt$$

and $G(x)$ is such that

$$x^{-\frac{1}{2}} F(x) = \sqrt{2} \int_0^\infty G_{2,4}^{2,1} \left(\frac{x^2 t^2}{4} \left| \begin{matrix} k-m-\frac{1}{4}-\frac{\nu}{2}, \frac{\nu}{2}-k+m+\frac{3}{4} \\ \frac{\nu}{2}+\frac{1}{4}, \frac{\nu}{2}+\frac{1}{4}+2m, \frac{1}{4}-\frac{\nu}{2}, \frac{1}{4}-\frac{\nu}{2}-2m \end{matrix} \right. \right) G(t) dt,$$

$$x^{-\frac{1}{2}} F\left(\frac{1}{x}\right) = \sqrt{2} \int_0^\infty G_{2,4}^{2,1} \left(\frac{x^2 t^2}{4} \left| \begin{matrix} k-m-\frac{1}{4}-\frac{\mu}{2}, \frac{\mu}{2}-k+m+\frac{3}{4} \\ \frac{\mu}{2}+\frac{1}{4}, \frac{\mu}{2}+\frac{1}{4}+2m, \frac{1}{4}-\frac{\mu}{2}, \frac{1}{4}-\frac{\mu}{2}-2m \end{matrix} \right. \right) G(t) dt,$$

then

$$g(x) = \int_0^\infty t^{-\frac{1}{2}} F(t) f(xt) dt$$

is $R_\mu(k, m)$, i.e. a solution of the integral equation

$$g(x) = \sqrt{2} \int_0^\infty G_{2,4}^{2,1} \left(\frac{x^2 t^2}{4} \left| \begin{matrix} k-m-\frac{1}{4}-\frac{\mu}{2}, \frac{\mu}{2}-k+m+\frac{3}{4} \\ \frac{\mu}{2}+\frac{1}{4}, \frac{\mu}{2}+\frac{1}{4}+2m, \frac{1}{4}-\frac{\mu}{2}, \frac{1}{4}-\frac{\mu}{2}-2m \end{matrix} \right. \right) g(t) dt,$$

where $G(x)$ is a continuous function in $(0, \infty)$ is $O(x^\lambda)$ for small x and $O(x^{-\delta})$ as $x \rightarrow \infty$, $f(x)$ belongs to $A(\alpha, a)$, $\text{Re}(\delta+1+2m-2k) > 0$, $\text{Re}(2\nu+\lambda+2+2m \pm 2m) > 0$, $\text{Re}(2\mu+\lambda+2+2m \pm 2m) > 0$ and $2m$ is not an integer or zero.

To justify the changes in the order of integration in (2.1) and (2.2), we see that u -integral is absolutely and uniformly convergent if $\text{Re} (2\nu + \lambda + 2 + 2m \pm 2m) > 0$, $\text{Re} (\delta + 1 + 2m - 2k) > 0$, $\text{Re} (2\mu + \lambda + 2 + 2m \pm 2m) > 0$, where $G(u) = O(u^\lambda)$ for small u and $O(x^{-\delta})$ as $u \rightarrow \infty$ and t -integral is absolutely and uniformly convergent if $f(t)$ belongs to $A(\alpha, a)$ and one of the repeated integrals is absolutely convergent if $f(t)$ belongs to $A(\alpha, a)$. Hence by de la Valle Poussin's theorem (Bromwich 1959, p. 504) the inversion of order of integration is justified. This proves the theorem. The argument also gives :

THEOREM 3. If $G(t)$ satisfies the same condition as in Theorem 2 then

$$g(x) = \int_0^\infty G(xu)f(u) du$$

is $R_\mu(k, m)$.

THEOREM 4. If $f(x)$ is $R_\mu(k, m)$ and $k(x)$ is $R_\nu(k, m)$, then the function

$$x^{-\frac{1}{2}}F(x) = \int_0^\infty k(y)f(xy) dy \quad \dots \quad (2.3)$$

is such that the $X_{\mu, k, m}$ -transform of $x^{-\frac{1}{2}}F(x)$ [$Q(x)$, say] is equal to $X_{\nu, k, m}$ -transform of $x^{-\frac{1}{2}}F\left(\frac{1}{x}\right)$.

PROOF: As $k(x)$ is $R_\nu(k, m)$, by known theorem (Narain 1956-57, p. 284)

$$x^{-\frac{1}{2}}F(x) = \frac{1}{2\pi i} \int_0^\infty f(xy) dy \int_{c-i\infty}^{c+i\infty} 2^{s/2} \frac{\Gamma_*\left(\frac{\nu}{2} + \frac{s}{2} + \frac{1}{4} + m \pm m\right)}{\Gamma\left(\frac{\nu}{2} + \frac{s}{2} + \frac{3}{4} + m - k\right)} \psi(s)y^{-s} ds,$$

where

$$\psi(s) = \psi(1-s).$$

Hence

$$x^{-\frac{1}{2}}F(x) = \frac{1}{2\pi i} \int_{c-i\infty}^{c+i\infty} 2^{s/2} \frac{\Gamma_*\left(\frac{\nu}{2} + \frac{s}{2} + \frac{1}{4} + m \pm m\right)}{\Gamma\left(\frac{\nu}{2} + \frac{s}{2} + \frac{3}{4} + m - k\right)} \psi(s) ds \int_0^\infty f(xy)y^{-s} dy.$$

The inversion of order of integration will be justified if $f(x)$ and $k(x)$ both belong to $A(\alpha, a)$.

Hence

$$x^{-\frac{1}{2}}F(x) = \frac{1}{2\pi i} \int_{c-i\infty}^{c+i\infty} 2^{s/2} \frac{\Gamma_*\left(\frac{\nu}{2} + \frac{s}{2} + \frac{1}{4} + m \pm m\right)}{\Gamma\left(\frac{\nu}{2} + \frac{s}{2} + \frac{3}{4} + m - k\right)} \psi(s)x^{s-1} ds \int_0^\infty u^{-s}f(u) du.$$

Now, as $f(x)$ is $R_\mu(k, m)$,

$$f(u) = \frac{1}{2\pi i} \int_{c'-i\infty}^{c'+i\infty} 2^{s/2} \frac{\Gamma_*\left(\frac{\mu}{2} + \frac{s}{2} + \frac{1}{4} + m \pm m\right)}{\Gamma\left(\frac{\mu}{2} + \frac{s}{2} + \frac{3}{4} + m - k\right)} \psi'(s) u^{-s} ds,$$

where

$$\psi'(s) = \psi'(1-s).$$

Hence, by Mellin's inversion formula,

$$\int_0^\infty u^{s-1} f(u) du = 2^{s/2} \frac{\Gamma_*\left(\frac{\mu}{2} + \frac{s}{2} + \frac{1}{4} + m \pm m\right)}{\Gamma\left(\frac{\mu}{2} + \frac{s}{2} + \frac{3}{4} + m - k\right)} \psi'(s),$$

changing s into $(1-s)$, we get

$$\int_0^\infty u^{-s} f(u) du = 2^{1-s} \frac{\Gamma_*\left(\frac{\mu}{2} + \frac{3}{4} - \frac{s}{2} + m \pm m\right)}{\Gamma\left(\frac{\mu}{2} + \frac{5}{4} - \frac{s}{2} + m - k\right)} \psi'(s).$$

Therefore

$$\begin{aligned} x^{-\frac{1}{2}}F(x) &= \frac{\sqrt{2}}{2\pi i} \int_{c-i\infty}^{c+i\infty} \frac{\Gamma_*\left(\frac{\nu}{2} + \frac{1}{4} + \frac{s}{2} + m \pm m\right) \Gamma_*\left(\frac{\mu}{2} + \frac{3}{4} - \frac{s}{2} + m \pm m\right)}{\Gamma\left(\frac{\nu}{2} + \frac{s}{2} + \frac{3}{4} + m - k\right) \Gamma\left(\frac{\mu}{2} + \frac{5}{4} - \frac{s}{2} + m - k\right)} \\ &\quad \times \psi(s)\psi'(s)x^{s-1} ds, \\ &= \frac{1}{2\pi i} \int_{1-c-i\infty}^{1-c+i\infty} \frac{\Gamma_*\left(\frac{\nu}{2} + \frac{3}{4} - \frac{s}{2} + m \pm m\right) \Gamma_*\left(\frac{\mu}{2} + \frac{1}{4} + \frac{s}{2} + m \pm m\right)}{\Gamma\left(\frac{\nu}{2} + \frac{5}{4} - \frac{s}{2} + m - k\right) \Gamma\left(\frac{\mu}{2} + \frac{3}{4} + \frac{s}{2} + m - k\right)} \\ &\quad \times \chi(s)x^{-s} ds, \quad \dots \dots \dots \quad (2.4) \end{aligned}$$

where

$$\chi(s) = \chi(1-s).$$

The form of this integral shows that $x^{-\frac{1}{2}}F(x)$ is a function such that $\chi_{\mu, k, m}$ -transform of $x^{-\frac{1}{2}}F(x)$ is equal to $\chi_{\nu, k, m}$ -transform of $x^{-\frac{1}{2}}F\left(\frac{1}{x}\right)$.

Putting $k+m = \frac{1}{2}$ in the above Theorem 4, we get a result in Hankel transform (Brij Mohan 1932).

THEOREM 5. If

$$Q(x) = \int_0^\infty \frac{1}{y} k(y) f\left(\frac{x}{y}\right) dy = \int_0^\infty \frac{1}{y} k\left(\frac{x}{y}\right) f(y) dy,$$

where $f(x)$ is $R_\mu(k, m)$ and $k(x)$ is $R_\nu(k, m)$ and both belong to $A(\alpha, a)$, then $Q(x)$ is of the form

$$Q(x) = \frac{1}{2\pi i} \int_{c-i\infty}^{c+i\infty} 2^{s/2} \frac{\Gamma_*\left(\frac{\mu}{2} + \frac{1}{4} + \frac{s}{2} + m \pm m\right) \Gamma_*\left(\frac{\nu}{2} + \frac{1}{4} + \frac{s}{2} + m \pm m\right)}{\Gamma\left(\frac{\mu}{2} + \frac{3}{4} + \frac{s}{2} + m - k\right) \Gamma\left(\frac{\nu}{2} + \frac{3}{4} + \frac{s}{2} + m - k\right)} \times \chi(s)x^{-s} ds, \quad \dots \dots \dots (2.5)$$

where

$$\chi(s) = \chi(1-s).$$

The theorem follows if we proceed in the same way as in Theorem 4.

As can be easily verified, $Q(x)$ is such that if $\chi_{\mu, k, m}$ -transform of $Q(x)$ is $x^{-1/2}F(x)$, then $\chi_{\nu, k, m}$ -transform of $Q(x)$ is $x^{-1/2}F\left(\frac{1}{x}\right)$.

On taking $k+m = \frac{1}{2}$ the above theorem will give a known result (Brij Mohan 1932).

§ 2.1. If, in the result of Theorem 4, we put $\mu = \nu$, we get by (2.4)

$$x^{-1/2}F(x) = \frac{1}{2\pi i} \int_{1-c-i\infty}^{1-c+i\infty} \lambda(s)x^{-s} ds, \quad \dots \dots \dots (2.6)$$

where

$$\lambda(s) = \lambda(1-s).$$

The symmetry of the result shows that we shall arrive at the same result, if we put

$$x^{-1/2}F(x) = \int_0^\infty k(xy)f(y) dy.$$

So we get a verification of the following theorem, which is a particular case of Parseval formula.

If $f(x)$ and $k(x)$ are $R_\nu(k, m)$, then

$$\int_0^\infty k(xy) f(y) dy = \int_0^\infty k(y) f(xy) dy. \quad \dots \dots (2.7)$$

§ 2.2. If $f(x)$ and $k(x)$ are both $R_\nu(k, m)$ and both belong to $A(\alpha, a)$, we have by (2.6), (2.3) and (2.7)

$$\begin{aligned} x^{-1/2}F(x) &= \int_0^\infty k(xy) f(y) dy \\ &= \frac{1}{2\pi i} \int_{1-c-i\infty}^{1-c+i\infty} \lambda(s) x^{-s} ds, \end{aligned}$$

where

$$\lambda(s) = \lambda(1-s).$$

Hence

$$\begin{aligned} F(x) &= \frac{1}{2\pi i} \int_{1-c-i\infty}^{1-c+i\infty} \lambda(s) x^{\frac{1}{2}-s} ds \\ &= \frac{1}{2\pi i} \int_{c-i\infty}^{c+i\infty} \lambda(1-s)x^{s-\frac{1}{2}} ds \\ &= \frac{1}{2\pi i} \int_{c-i\infty}^{c+i\infty} \lambda(s)x^{s-\frac{1}{2}} ds = F\left(\frac{1}{x}\right). \end{aligned}$$

§ 3. *Examples.*—Using the results of §2.1, we can derive some examples of self-reciprocal function, i.e. solution of the integral equation.

$$f(x) = \sqrt{2} \int_0^\infty G_{2,4}^{2,1} \left(\frac{x^2 y^2}{4} \left| \begin{matrix} k-m-\frac{1}{4}-\frac{\nu}{2}, \frac{\nu}{2}-k+m+\frac{3}{4} \\ \frac{\nu}{2}+\frac{1}{4}, \frac{\nu}{2}+\frac{1}{4}+2m, \frac{1}{4}-\frac{\nu}{2}, \frac{1}{4}-\frac{\nu}{2}-2m \end{matrix} \right. \right) f(y) dy.$$

(i) Let us take (Narain 1956-57)

$$\begin{aligned} k(x) &= G_{1,2}^{2,0} \left(\frac{x^2}{2} \left| \begin{matrix} \frac{\nu}{2}+\frac{3}{4}+m-k \\ \frac{\nu}{2}+\frac{1}{4}, \frac{\nu}{2}+\frac{1}{4}+2m \end{matrix} \right. \right) \\ &= x^{\nu+2m-\frac{1}{2}} e^{-\frac{1}{2}x^2} W_{k,m} \left(\frac{x^2}{2} \right), \end{aligned}$$

then

$$F(x) = \sqrt{x} \int_0^\infty G_{1,2}^{2,0} \left(\frac{x^2}{2} \left| \begin{matrix} \frac{\nu}{2}+\frac{3}{4}+m-k \\ \frac{\nu}{2}+\frac{1}{4}, \frac{\nu}{2}+\frac{1}{4}+2m \end{matrix} \right. \right) f(y) dy.$$

Inverting this integral by Inversion formula of Meijer-Laplace transform (Bhise 1959) we get a second solution.

(ii) Let us have (Saxena unpublished)

$$k(x) = G_{2,4}^{2,1} \left(\frac{x^4}{16} \left| \begin{matrix} -\frac{\nu}{4}-\frac{1}{8}-\frac{m}{2}+\frac{k}{2}, \frac{\nu}{4}+\frac{1}{8}+\frac{m}{2}-\frac{k}{2} \\ \frac{\nu}{4}+\frac{1}{8}, \frac{\nu}{4}+\frac{1}{8}+m, \frac{1}{8}-\frac{\nu}{4}, \frac{1}{8}-\frac{\nu}{4}-m \end{matrix} \right. \right),$$

then

$$F(x) = \sqrt{x} \int_0^\infty G_{2,4}^{2,1} \left(\frac{x^4 y^4}{16} \left| \begin{matrix} -\frac{\nu}{4}-\frac{1}{8}-\frac{m}{2}+\frac{k}{2}, \frac{\nu}{4}+\frac{1}{8}+\frac{m}{2}-\frac{k}{2} \\ \frac{\nu}{4}+\frac{1}{8}, \frac{\nu}{4}+\frac{1}{8}+m, \frac{1}{8}-\frac{\nu}{4}, \frac{1}{8}-\frac{\nu}{4}-m \end{matrix} \right. \right) f(y) dy.$$

Now using the inversion formula (Fox 1961) for the transform (1.3), we can get the third solution.

§ 3.1. We work out the last solution in a slightly different form, thus:

Let us take

$$\psi(\beta) = \beta^{\frac{1}{2}} \int_0^{\infty} G_{2,4}^{2,1} \left(\frac{x^4 \beta^2}{16} \left| \begin{matrix} -\frac{\nu}{4} - \frac{1}{8} - \frac{m}{2} + \frac{k}{2}, \frac{\nu}{4} + \frac{3}{8} + \frac{m}{2} - \frac{k}{2} \\ \frac{\nu}{4} + \frac{1}{8}, \frac{\nu}{4} + \frac{1}{8} + m, \frac{1}{8} - \frac{\nu}{4}, \frac{1}{8} - \frac{\nu}{4} - m \end{matrix} \right. \right) f(x) dx.$$

If $f(x)$ is $R_{\nu}(k, m)$, then

$$f(x) = \sqrt{2} \int_0^{\infty} G_{2,4}^{2,1} \left(\frac{x^2 y^2}{4} \left| \begin{matrix} k - m - \frac{1}{4} - \frac{\nu}{2}, \frac{\nu}{2} + \frac{3}{4} - k + m \\ \frac{\nu}{2} + \frac{1}{4}, \frac{\nu}{2} + \frac{1}{4} + 2m, \frac{1}{4} - \frac{\nu}{2}, \frac{1}{4} - \frac{\nu}{2} - 2m \end{matrix} \right. \right) f(y) dy.$$

Multiplying both the sides by

$$\beta^{\frac{1}{2}} G_{2,4}^{2,1} \left(\frac{x^4 \beta^2}{16} \left| \begin{matrix} -\frac{\nu}{4} - \frac{1}{8} - \frac{m}{2} + \frac{k}{2}, \frac{\nu}{4} + \frac{3}{8} + \frac{m}{2} - \frac{k}{2} \\ \frac{\nu}{4} + \frac{1}{8}, \frac{\nu}{4} + \frac{1}{8} + m, \frac{1}{8} - \frac{\nu}{4}, \frac{1}{8} - \frac{\nu}{4} - m \end{matrix} \right. \right)$$

and integrating, we get

$$\begin{aligned} \psi(\beta) &= \sqrt{2} \beta^{\frac{1}{2}} \int_0^{\infty} G_{2,4}^{2,1} \left(\frac{x^4 \beta^2}{16} \left| \begin{matrix} -\frac{\nu}{4} - \frac{1}{8} - \frac{m}{2} + \frac{k}{2}, \frac{\nu}{4} + \frac{3}{8} + \frac{m}{2} - \frac{k}{2} \\ \frac{\nu}{4} + \frac{1}{8}, \frac{\nu}{4} + \frac{1}{8} + m, \frac{1}{8} - \frac{\nu}{4}, \frac{1}{8} - \frac{\nu}{4} - m \end{matrix} \right. \right) dx \\ &\quad \times \int_0^{\infty} G_{2,4}^{2,1} \left(\frac{x^2 y^2}{4} \left| \begin{matrix} k - m - \frac{1}{4} - \frac{\nu}{2}, \frac{\nu}{2} + \frac{3}{4} - k + m \\ \frac{\nu}{2} + \frac{1}{4}, \frac{\nu}{2} + \frac{1}{4} + 2m, \frac{1}{4} - \frac{\nu}{2}, \frac{1}{4} - \frac{\nu}{2} - 2m \end{matrix} \right. \right) f(y) dy \\ &= \sqrt{2} \beta^{\frac{1}{2}} \int_0^{\infty} f(y) dy \int_0^{\infty} G_{2,4}^{2,1} \left(\frac{x^2 y^2}{4} \left| \begin{matrix} k - m - \frac{1}{4} - \frac{\nu}{2}, \frac{\nu}{2} + \frac{3}{4} - k + m \\ \frac{\nu}{2} + \frac{1}{4}, \frac{\nu}{2} + \frac{1}{4} + 2m, \frac{1}{4} - \frac{\nu}{2}, \frac{1}{4} - \frac{\nu}{2} - 2m \end{matrix} \right. \right) \\ &\quad \times G_{2,4}^{2,1} \left(\frac{x^4 \beta^2}{16} \left| \begin{matrix} -\frac{\nu}{4} - \frac{1}{8} - \frac{m}{2} + \frac{k}{2}, \frac{\nu}{4} + \frac{3}{8} + \frac{m}{2} - \frac{k}{2} \\ \frac{\nu}{4} + \frac{1}{8}, \frac{\nu}{4} + \frac{1}{8} + m, \frac{1}{8} - \frac{\nu}{4}, \frac{1}{8} - \frac{\nu}{4} - m \end{matrix} \right. \right) dx, \end{aligned}$$

replacing x^2 by x and using (Saxena 1960), (Erdélyi *et al.* 1953) we have

$$\begin{aligned} \psi(\beta) &= \beta^{-\frac{1}{2}} \int_0^{\infty} G_{2,4}^{2,1} \left(\frac{y^4}{16 \beta^2} \left| \begin{matrix} -\frac{\nu}{4} - \frac{1}{8} + \frac{k}{2} - \frac{m}{2}, \frac{\nu}{4} + \frac{1}{8} + \frac{m}{2} - \frac{k}{2} \\ \frac{\nu}{4} + \frac{1}{8}, \frac{\nu}{4} + \frac{1}{8} + m, \frac{1}{8} - \frac{\nu}{4}, \frac{1}{8} - \frac{\nu}{4} - m \end{matrix} \right. \right) f(y) dy \\ &= \psi \left(\frac{1}{\beta} \right). \end{aligned}$$

Now

$$\begin{aligned}\psi(\beta) &= \beta^{\frac{1}{2}} \int_0^{\infty} G_{2,4}^{2,1} \left(\frac{x^4 \beta^2}{16} \left| \begin{array}{c} -\frac{\nu}{4} - \frac{1}{8} + \frac{k}{2} - \frac{m}{2}, \frac{\nu}{4} + \frac{3}{8} + \frac{m}{2} - \frac{k}{2} \\ \frac{\nu}{4} + \frac{1}{8}, \frac{\nu}{4} + \frac{1}{8} + m, \frac{1}{8} - \frac{\nu}{4}, \frac{1}{8} - \frac{\nu}{4} - m \end{array} \right. \right) f(x) dx \\ &= \beta^{\frac{1}{2}} \int_0^{\infty} G_{2,4}^{2,1} \left(\frac{\xi^2 \beta^2}{4} \left| \begin{array}{c} -\frac{\nu}{4} - \frac{1}{8} + \frac{k}{2} - \frac{m}{2}, \frac{\nu}{4} + \frac{3}{8} + \frac{m}{2} - \frac{k}{2} \\ \frac{\nu}{4} + \frac{1}{8}, \frac{\nu}{4} + \frac{1}{8} + m, \frac{1}{8} - \frac{\nu}{4}, \frac{1}{8} - \frac{\nu}{4} - m \end{array} \right. \right) \frac{f(2\sqrt{\xi})}{\sqrt{\xi}} d\xi.\end{aligned}$$

Therefore by inversion formula (Fox 1961), we get

$$\xi^{-\frac{1}{2}} f(2\sqrt{\xi}) = \int_0^{\infty} G_{2,4}^{2,1} \left(\frac{\xi^2 t^2}{4} \left| \begin{array}{c} -\frac{\nu}{4} - \frac{1}{8} + \frac{k}{2} - \frac{m}{2}, \frac{\nu}{4} + \frac{3}{8} + \frac{m}{2} - \frac{k}{2} \\ \frac{\nu}{4} + \frac{1}{8}, \frac{\nu}{4} + \frac{1}{8} + m, \frac{1}{8} - \frac{\nu}{4}, \frac{1}{8} - \frac{\nu}{4} - m \end{array} \right. \right) \psi(t) dt.$$

Thus

$$\frac{2}{x} f(x) = \int_0^{\infty} G_{2,4}^{2,1} \left(\frac{x^4 t^2}{16} \left| \begin{array}{c} -\frac{\nu}{4} - \frac{1}{8} + \frac{k}{2} - \frac{m}{2}, \frac{\nu}{4} + \frac{3}{8} + \frac{m}{2} - \frac{k}{2} \\ \frac{\nu}{4} + \frac{1}{8}, \frac{\nu}{4} + \frac{1}{8} + m, \frac{1}{8} - \frac{\nu}{4}, \frac{1}{8} - \frac{\nu}{4} - m \end{array} \right. \right) \psi(t) dt.$$

Therefore, we obtain

$$f(x) = \frac{x}{2} \int_0^{\infty} G_{2,4}^{2,1} \left(\frac{x^4 t^2}{16} \left| \begin{array}{c} -\frac{\nu}{4} - \frac{1}{8} + \frac{k}{2} - \frac{m}{2}, \frac{\nu}{4} + \frac{3}{8} + \frac{m}{2} - \frac{k}{2} \\ \frac{\nu}{4} + \frac{1}{8}, \frac{\nu}{4} + \frac{1}{8} + m, \frac{1}{8} - \frac{\nu}{4}, \frac{1}{8} - \frac{\nu}{4} - m \end{array} \right. \right) \psi(t) dt$$

where

$$\psi(t) = \psi\left(\frac{1}{t}\right).$$

This is the third solution.

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