

# COLLINEATION AND MOTION IN RELATIVISTIC THERMODYNAMICS

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In this paper it has been shown that Ricci collineation along flow vector implies motion in the space-time of relativistic thermodynamics. The theory of relativistic thermodynamics used here is Cattaneo's theory.

## INTRODUCTION

Let  $g_{ab}$  denote the metric tensor in the 4-dimensional space-time  $V_4$  of relativistic thermodynamics with the signature  $(- - - +)$ . We will represent partial derivative by a comma and covariant derivative by a semicolon. Roman indices  $a, b, \dots$  will take values 1 to 4. Since metric tensor is covariant constant,

$$g_{ab;c} = 0.$$

Following Cattaneo's theory the stress energy tensor for the relativistic thermodynamical fluid is given by

$$T_{ab} = (\rho + p) u_a u_b - p g_{ab} - Q_a u_b - Q_b u_a, \quad \dots(1)$$

where  $\rho$  is the total energy density of the fluid,  $p$  is the pressure, and  $Q_a$  are the components of the heat flux vector and  $U^a$  is the time like flow vector such that

$$Q_a u^a = 0 \quad \dots(2a)$$

$$u_a u^a = 1 \quad \dots(2b)$$

$$Q_a Q^a = -Q^2 \quad \dots(2c)$$

Since for the thermodynamical fluid stress tensor is given by Eq. (1), Ricci tensor is obtained from Eq. (1) in the form

$$R_{ab} = \psi (\rho + p) u_a u_b - \frac{\psi}{2} (\rho - p) g_{ab} - Q_a u_b - Q_b u_a$$

$$R_{ab} = A u_a u_b - B g_{ab} - \lambda Q_a u_b - \lambda Q_b u_a \quad \dots(3)$$

$$A = \psi (\rho + p) \quad \dots(4)$$

$$B = \frac{\psi}{2} (\rho - p). \quad \dots(5)$$

Equations (4) and (5) yield

$$A - B = \frac{\psi}{2} (\rho + 3p). \quad \dots(6)$$

In consequence of Eqs. (2) and (3) we get

$$R_{ab} u^a = (A - B) u_b - Q_b$$

$$R_{ab} u^b u^b = (A - B)$$

$$R_{ab} Q^a = -BQ_b + Q^2 u_b$$

$$R_{ab} Q^a Q^b = BQ^2.$$

A vector  $Q^a$  is said to be harmonic (Nordvedt & Pagel 1962) if

$$Q_{a;b} - Q_{ab} = 0 \tag{7}$$

$$Q_{;a}^a = 0$$

RICCI COLLINEATION

Katzine *et al.* (1969) defined Ricci collineation as the point transformation

$$X^i \rightarrow X^i + \xi^i \delta t,$$

which keeps the form of the Ricci tensor invariant along the field vector  $\xi^i$ , that is for Ricci collineation

$$\mathcal{L}_\xi R_a = 0 \text{ when } \mathcal{L} \text{ denotes the Lie derivative.}$$

When a space-time admits groups of motions it is obvious that it admits the corresponding Ricci collineation

$$\text{i.e., } \frac{\mathcal{L}}{k} g_{ab} = 0 \Rightarrow \frac{\mathcal{L}}{k} R_{ab} = 0.$$

However, even if the converse is not true, in general. We have shown that in this case converse is also true under certain assumptions.

*Theorem 1 — In relativistic thermodynamical non-Zeldovich fluid if  $Q^a$  is harmonic and  $p_{;k} u^k = 0$ , then*

$$\frac{\mathcal{L}}{u} R_{ab} = 0 \text{ implies that the world lines of material particles are geodesic.}$$

$$\text{PROOF : } \frac{\mathcal{L}}{u} R_{ab} = R_{ab;k} u^k + R_{kb} u^k_{;a} + R_{ak} u^k_{;b} \tag{8}$$

With the help of Eq. (3) and the assumption that the heat flux vector is harmonic i.e.,

$$Q_{a;b} - Q_{b;a} = 0$$

Eg. (8) reduces to

$$\frac{\mathcal{L}}{u} R_{ab} = A_{;k} u_a u_b u^k - B_{;k} g_{ab} u^k + \dot{u}_a (A u_b = \chi Q_b)$$

$$+ \dot{u}_b (A u_a - \xi \chi Q_a) - B (u_{a;b} + u_{b;a}) \tag{9}$$

where  $u_a \stackrel{\text{def}}{=} u_{a;k} u^k$ .

$$u^a u^b \frac{\mathcal{L}}{u} R_{ab} = (A - B)_{;k} u^k$$

$$\frac{\mathfrak{L}}{u} R_{ab} = 0 \Rightarrow u^a u^b \frac{\mathfrak{L}}{u} R_{ab} = 0.$$

Hence in consequence of Eq. (6)

$$(A - B)_{;k} u^k = 0 \Rightarrow (\rho + 3p)_{;k} u^k = 0.$$

The hypothesis  $p_{;k} u^k = 0$   
 $\rho_{;k} u^k = 0.$

..(10)

Also  $u^a \frac{\mathfrak{L}}{u} R_{ab} = 0 \Rightarrow$

$$(A - B)_{;k} u^k u_b + (A - B) u_b = 0.$$

..(11)

Using  $(A - B)_{;k} u^k = 0$  in Eq. (11) we obtain  $(A - B) \dot{u}_b = 0.$

But since

$$A - B \neq 0,$$

$$U_b = 0,$$

..(12)

which is the required condition.

*Theorem 2 — In relativistic thermodynamical non-Zeldovich fluid if  $Q^a$  is harmonic and  $p_{;k} u^k = 0$ , then*

$$\frac{\mathfrak{L}}{u} R_{ab} = 0 \Rightarrow \frac{\mathfrak{L}}{u} g_{ab} = 0.$$

PROOF : We have

$$\begin{aligned} \frac{\mathfrak{L}}{u} R_{ab} &= R_{ab;k} u^k + R_{kb} u^k_{;a} + R_{ak} u^k_{;b} \\ &= [A u_a u_b - B g_{ab} - Q_a u_b - Q_b u_a]_{;k} u^k \\ &\quad + [A u_k u - B g_{kb} - Q_k u_b - Q_b u_k] u^k_{;a} + [A u_a u_k - B g_{ak} \\ &\quad - Q_a u_k - Q_k u_a] u^k_{;b}. \end{aligned}$$

Since  $Q^a$  is harmonic this equation reduces to

$$\frac{\mathfrak{L}}{u} R_{ab} = T_1 + T_2 - T_3$$

where

$$\begin{aligned} T_1 &\underline{\text{def}} A_{;k} u_a u_b u^k - B_{;k} g_{ab} u^k \\ T_2 &\underline{\text{def}} u_a (A u_b - Q_b) + (A u_a - Q_a) \dot{u}_b \\ T_3 &\underline{\text{def}} B (u_{a;b} + u_{b;a}). \end{aligned}$$

In order to prove  $\frac{\xi}{u} R_{ab} = 0 \Rightarrow \frac{\xi}{u} g_{ab} = 0$ , we have to prove that  $T_1$  and  $T_2$  vanish. Eqs. (10) and (12) show that  $T_1$  and  $T_2$  will vanish.

$$\frac{\xi}{u} R_{ab} = 0, \quad T_1 = 0, \quad T_2 = 0$$

together imply that

$$B(\lambda u_{a;b} + \lambda u_{b;a}) = 0.$$

Since the fluid is non-zeldovich that is  $\rho \neq p$ , therefore

$$B \neq 0 \Rightarrow U_{a;b} + U_{b;a} = 0$$

which is the required result.

*Remark :* In relativistic thermodynamical non-zeldovich fluid for harmonic heat flux vector and uniform pressure Ricci-collineation along flow vector implies that the world lines of material particles are geodesic. Ricci collineation implies motion along flow vector in the four dimensional space time  $V_4$  with the above mentioned assumptions.

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