

**DEFLECTION OF AN ELASTIC CIRCULAR PLATE WITH CLAMPED EDGE
RESTING ON A VISCOELASTIC FOUNDATION AND HAVING A
CONCENTRATED LOAD AT THE CENTRE**

by **MADAN MOHAN GARAI**, *Jadavpur University, Calcutta 32*

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In the present paper the complete solution of the problem of the deflection of an elastic circular plate with fixed edge resting on a viscoelastic foundation has been obtained with the assumption that there is a concentrated normal load at the centre.

1. INTRODUCTION

Following Reissner (1958) we take the equation of motion of an elastic plate on viscoelastic foundation as

$$D \nabla_1^4 w + K \left(1 + \lambda \frac{\partial}{\partial t} \right) w + \rho h \frac{\partial^2 w}{\partial t^2} = q, \quad \dots(1.1)$$

where

w = transverse displacement of the plate,

D = bending stiffness factor of the plate,

$K \left(1 + \lambda \frac{\partial}{\partial t} \right)$ = modulus of viscoelastic foundation
(K and λ being constants),

ρ = uniform density of the plate,

h = its thickness, and

q = transverse load per unit area.

For clamped edge, boundary conditions for a circular plate of radius a are

$$\left. \begin{array}{l} \text{(i) } (w)_{r=a} = 0 \\ \text{(ii) } \left(\frac{\partial w}{\partial r} \right)_{r=a} = 0. \end{array} \right\} \quad \dots(1.2)$$

2. SOLUTION OF THE PROBLEM

To solve eqn. (1.1) we shall first consider the solution of the equation

$$\nabla_1^4 W = k^4 W \quad (k \text{ is a constant}). \quad \dots(2.1)$$

For symmetrical solution, the above equation expressed in polar co-ordinates can be written as

$$\left[\left(\frac{d^2}{dr^2} + \frac{1}{r} \cdot \frac{d}{dr} \right)^2 - k^4 \right] W = 0 \quad \dots(2.2)$$

The solution of this equation suitable for a complete disc is

$$W = A_1 J_0(kr) + B_1 I_0(kr) \quad \dots(2.3)$$

in which J_0 is the Bessel function and I_0 is the modified Bessel function of zero order and of first kind.

The boundary condition (i) of (1.2) gives

$$W = I_0(ka) J_0(kr) - J_0(ka) I_0(kr) = Z_0(kr) \text{ (say)} \quad \dots(2.4)$$

The second boundary condition

$$\left(\frac{\partial w}{\partial r} \right)_{r=a} = 0$$

i.e. $Z'_0(ka) = 0$

or

$$\frac{J_1(ka)}{J_0(ka)} + \frac{I_1(ka)}{I_0(ka)} = 0 \quad \dots(2.5)$$

If $k_1, k_2, k_3, \dots, k_n, \dots$ are the roots of this equation, then $W = Z_0(k_1 r), Z_0(k_2 r), Z_0(k_3 r), \dots, Z_0(k_n r), \dots$ are solutions of the equation (2.1) which satisfy the boundary conditions.

Let us put

$$w = \sum_{n=1}^{\infty} A_n(t) Z_0(k_n r) \quad \dots(2.6)$$

Again if $q = P\delta(r)/2\pi r$ where P is a constant, and $\delta(r)$ is Dirac delta function, we can write

$$P\delta(r)/2\pi r = \sum_{n=1}^{\infty} B_n Z_0(k_n r) \quad \dots(2.7)$$

B_n being, a constant.

It can be readily shown that the function $Z_0(k_n r)$ are orthogonal with the following properties similar to those for Bessel functions :

$$\int_0^a r Z_0(k_n r) Z_0(k_m r) dr \simeq 0$$

$$\int_0^a r Z_0^2(k_n r) dr = \frac{1}{2} a^2 [I_0^2(k_n a) J_1^2(k_n a) - J_0^2(k_n a) I_1^2(k_n a) + 2 J_0^2(k_n a) I_0^2(k_n a)]. \quad \dots(2.8)$$

Multiplying both sides of (2.7) by $r Z_0(k_n r) dr$, we obtain from (2.7),

$$\frac{P}{2\pi} \int_0^a \delta(r) Z_0(k_n r) dr = B_n \int_0^a r Z_0^2(k_n r) dr$$

or

$$B_n = \frac{P}{2\pi} \cdot \frac{[I_0(k_n a) - J_0(k_n a)]}{\frac{a^2}{2} [I_0^2(k_n a) J_1^2(k_n a) - J_0^2(k_n a) I_1^2(k_n a) + 2 J_0^2(k_n a) I_0^2(k_n a)]}. \quad \dots(2.9)$$

Putting

$$w = \sum_{n=1}^{\infty} A_n(t) Z_0(k_n r) \text{ and writing } q = \sum_{n=1}^{\infty} B_n Z_0(k_n r)$$

we obtain eqn. (1.1) as

$$D \sum_{n=1}^{\infty} A_n(t) k_n^4 Z_0(k_n r) + K \left(1 + \lambda \frac{\partial}{\partial t} \right) \sum_{n=1}^{\infty} A_n(t) Z_0(k_n r) + \rho h \sum_{n=1}^{\infty} \frac{d^2 A_n(t)}{dt^2} Z_0(k_n r) = \sum_{n=1}^{\infty} B_n Z_0(k_n r).$$

This will be satisfied if for all values of $\eta = 1, 2, 3, \dots$,

$$D A_n(t) k_n^4 + K \left(1 + \lambda \frac{\partial}{\partial t} \right) A_n(t) + \rho h \frac{d^2 A_n(t)}{dt^2} = B_n.$$

Taking Laplace transform and using the symbol

$$\bar{\varphi}(p) = \int_0^{\infty} e^{-pt} \varphi(t) dt$$

we have

$$[(Dk_n^4 + K) + K\lambda p + \rho h p^2] \bar{A}_n(p) = \frac{B_n}{p}.$$

It is assumed that $w = \frac{\partial w}{\partial t} = 0$, when $t = 0$.

$$\begin{aligned} \therefore \bar{A}_n(p) &= \frac{B_n}{p [\rho h p^2 + K \lambda p + (Dk_n^4 + K)]} \\ &= B_n \left[\frac{1}{(Dk_n^4 + K) p} - \frac{p \rho h}{(Dk_n^4 + K) \{\rho h p^2 + K \lambda p + (Dk_n^4 + K)\}} \right. \\ &\quad \left. - \frac{K \lambda}{(Dk_n^4 + K) \{\rho h p^2 + K \lambda p + (Dk_n^4 + K)\}} \right] \\ &= B_n \left[\frac{1}{(Dk_n^4 + K) p} - \frac{1}{Dk_n^4 + K} \cdot \frac{p + \frac{K \lambda}{2 \rho h}}{\left(p + \frac{K \lambda}{2 \rho h} \right)^2 + \left(\frac{Dk_n^4 + K}{\rho h} - \frac{K^2 \lambda^2}{4 \rho^2 h^2} \right)} \right. \\ &\quad \left. - \frac{1}{Dk_n^4 + K} \cdot \frac{K \lambda}{2 \rho h} \cdot \frac{1}{\left(p + \frac{K \lambda}{2 \rho h} \right)^2 + \left(\frac{Dk_n^4 + K}{\rho h} - \frac{K^2 \lambda^2}{4 \rho^2 h^2} \right)} \right]. \end{aligned}$$

Taking inverse transform we get

$$\begin{aligned} A_n(t) &= B_n \left[\frac{1}{Dk_n^4 + K} - \frac{1}{Dk_n^4 + K} e^{-(K \lambda / 2 \rho h) t} \cos \sqrt{\frac{Dk_n^4 + K}{\rho h} - \frac{K^2 \lambda^2}{4 \rho^2 h^2}} t \right. \\ &\quad \left. - \frac{1}{Dk_n^4 + K} \cdot \frac{K \lambda}{2 \rho h} e^{-(K \lambda / 2 \rho h) t} \sin \sqrt{\frac{Dk_n^4 + K}{\rho h} - \frac{K^2 \lambda^2}{4 \rho^2 h^2}} t \right] \end{aligned}$$

$$\begin{aligned} \therefore w &= \sum_{n=1}^{\infty} B_n Z_0(k_n r) \left[\frac{1}{Dk_n^4 + K} - \frac{1}{Dk_n^4 + K} e^{-(K \lambda / 2 \rho h) t} \right. \\ &\quad \times \cos \sqrt{\frac{Dk_n^4 + K}{\rho h} - \frac{K^2 \lambda^2}{4 \rho^2 h^2}} t - \frac{K \lambda}{\rho h (Dk_n^4 + K)} \\ &\quad \left. \times e^{-(K \lambda / 2 \rho h) t} \sin \sqrt{\frac{Dk_n^4 + K}{\rho h} - \frac{K^2 \lambda^2}{4 \rho^2 h^2}} t \right]. \end{aligned}$$

Roots of eqn. (2.5) are

$$\begin{aligned} k_1 a &= 3.196, k_2 a = 6.306, k_3 a = 9.440, \\ k_4 a &= 12.577, k_5 a = 15.716, \text{ etc.} \end{aligned} \quad \dots(2.10)$$

For large t i.e. when $t \rightarrow \infty$

$$\begin{aligned} w &= \sum_{n=1}^{\infty} B_n Z_0(k_n r) \frac{1}{Dk_n^4 + K} \\ &= \frac{P}{2\pi} \sum_{n=1}^{\infty} \frac{I_0(k_n a) - J_0(k_n a)}{\frac{a^2}{2} [I_0^2(k_n a) J_1^2(k_n a) - J_0^2(k_n a) I_1^2(k_n a) + 2 J_0^2(k_n a) I_0^2(k_n a)]} \\ &\quad \times \frac{Z_0(k_n r)}{Dk_n^4 + K}. \end{aligned} \quad \dots(2.11)$$

Since

$$[Z_0(k_n r)]_{r=0} = I_0(k_n a) - J_0(k_n a) \quad [\text{vide (2.4)}]$$

we have from (2.10)

$$w = \frac{P}{\pi} \cdot \frac{5.4 a^4}{a^2(D 104.8 + K a^4)} \quad (\text{approx.}).$$

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